Hosed vs. unhosed: global response to interruptions of the Atlantic Meridional Overturning, with and without freshwater forcing

N. Brown and E. D. Galbraith

Dept. of Earth and Planetary Science, McGill University, Montreal QC H3A 2A7, Canada

Received: 29 August 2015 – Accepted: 5 September 2015 – Published: 5 October 2015

Correspondence to: E. D. Galbraith (eric.d.galbraith@gmail.com)

Published by Copernicus Publications on behalf of the European Geosciences Union.
Abstract

It is well known that glacial periods were punctuated by abrupt climate changes, with large impacts on air temperature, precipitation, and ocean circulation across the globe. However, the long-held idea that freshwater forcing, caused by massive iceberg discharges, was the driving force behind these changes has been questioned in recent years. This throws into doubt the abundant literature on modelling abrupt climate change through “hosing” experiments, whereby the Atlantic Meridional Overturning Circulation (AMOC) is interrupted by an injection of freshwater to the North Atlantic: if some, or all, abrupt climate change was not driven by freshwater input, could its character have been very different than the typical hosed experiments? Here, we take advantage of a global coupled ocean–atmosphere model that exhibits spontaneous, unhosed oscillations in AMOC strength, in order to examine how the global imprint of AMOC variations depends on whether or not it is the result of external freshwater input. The results imply that, to first order, the ocean–ice–atmosphere dynamics associated with an AMOC weakening dominate the global response, regardless of whether or not freshwater input is the cause. The exception lies in the impact freshwater inputs can have on the strength of other polar haloclines, particularly the Southern Ocean, to which freshwater can be transported relatively quickly after injection in the North Atlantic.

1 Introduction

“Abrupt” climate changes were initially identified as decadal-centennial temperature changes in Greenland ice deposited during the last ice age, and subsequently recognized as globally-coherent climate shifts (Voelker, 2002; Alley et al., 2003). Abrupt climate changes involved changes in the latitudinal extent and strength of the AMOC (Clark et al., 2002), but also impacted global patterns of precipitation, surface air temperature, ocean biogeochemistry and atmospheric trace gas composition (Clement and
Abrupt climate changes were associated, early on, with dramatic layers of ice-rafted detritus that blanketed the North Atlantic during brief intervals of the last ice age (Heinrich, 1988), leading to the idea that recurring pulses of freshwater input had been the cause of AMOC interruptions. As evocatively described by Broecker (1994), armadas of icebergs, periodically discharged from the northern ice sheets to melt across the North Atlantic (MacAyeal, 1993), would have spread a freshwater cap that impeded convection and consequently, through the Stommel (1961) feedback, would have thrown a wrench in the overturning. Inspired by this idea, generations of numerical models have been subjected to freshwater “hosing” experiments, whereby the sensitivity of the AMOC to varying degrees of freshwater input have been plumbed, and the responses shown to vary as a function of background climate state and experimental design (Fanning and Weaver, 1997; Ganopolski and Rahmstorf, 2001; Rahmstorf et al., 2005; Stouffer et al., 2006; Krebs and Timmermann, 2007; Hu et al., 2008; Otto-Bliesner and Brady, 2010; Kageyama et al., 2013; Gong et al., 2013; Roberts et al., 2014). Multiple studies have shown a good degree of consistency between aspects of these hosing simulations and the observed global signatures of abrupt climate change (Liu et al., 2009; Menviel et al., 2014).

However, other model simulations have shown that spontaneous changes in the AMOC can occur in the absence of freshwater inputs (Winton, 1993; Sakai and Peltier, 1997; Hall and Stouffer, 2001; Ganopolski and Rahmstorf, 2002; Schulz, 2002; Loving and Vallis, 2005; Wang and Mysak, 2006; Friedrich et al., 2010; Kim et al., 2012; Drijfhout et al., 2013; Peltier and Vettoretti, 2014; Vettoretti and Peltier, 2015). Although uncommon, these “unhosed” simulations show that the AMOC can vary as a result of processes internal to the ocean–atmosphere system, most importantly through the existence of multiple stable states of the sea ice extent in the North Atlantic (Li, 2005; Li et al., 2010; Siddall et al., 2010; Petersen et al., 2013). In addition, there are features of the observational records that are inconsistent with iceberg armadas having been the driving force behind all episodes of abrupt climate change (Marshall...
and Koutnik, 2006). Periods of rapid Greenland cooling are typically divided between “Heinrich events”, for which widespread ice-rafted detritus is found (Hemming, 2004), and “Dansgaard-Oeschger stadials” (Dansgaard et al., 1993), which include all abrupt Greenland coolings, but are not associated with widespread ice rafted detritus. In Greenland ice core records, Heinrich events are very similar to Dansgaard-Oeschger stadials, despite the apparent contrast in the associated amount of ice rafted detritus and, presumably, the consequent freshwater input.

What’s more, the arrival of ice-rafted detritus does not seem to precede changes in the AMOC, where the two are recorded together. An analysis of Heinrich event 2 in the NW Atlantic showed that weakening of the AMOC preceded the widespread deposition of ice-rafted detritus by approximately 2 ky (Gutjahr and Lippold, 2011), while a statistical analysis of the temporal relationship between ice-rafted detritus near Iceland and dozens of cold events in Greenland suggested that ice-rafted detritus deposition generally lags the onset of cooling by a significant amount (Barker et al., 2015). These observations throw further doubt on the role of iceberg melting as the universal driver of abrupt climate change. In fact, it may be more likely that iceberg release is more often a consequence of AMOC interruptions, rather than necessarily being their cause, as raised by the identifications of subsurface warming during AMOC interruptions in hosed model simulations (Mignot et al., 2007). Such intermediate-depth warming of the North Atlantic, resulting from the reduced release of oceanic heat at high latitudes, would have melted floating ice shelves at their bases, contributing to ice sheet collapse (Shaffer, 2004; Flückiger et al., 2006; Alvarez-Solas et al., 2010; Marcott et al., 2011). Thus, although there is good evidence that large iceberg armadas were released during some stadials, and would have freshened the North Atlantic accordingly, they were not necessarily the main causal factor involved in all stadials. This raises the question: does the global footprint of an AMOC interruption depend on its cause, or is any AMOC interruption the same, regardless of whether or not freshwater forcing was behind it? How does the importance of the causal driver compare with the sensitivity to the background climate state?
In order to explore these questions, we make use of a large number of long water-hosing simulations with CM2Mc, a state-of-the-art Earth System model, to show how the global response to hosing varies between a preindustrial and glacial background state, and under different orbital forcings. In addition, we take advantage of the fact that the same model exhibits previously undescribed spontaneous AMOC interruptions and resumptions, which appear very similar to stadial-interstadial variability, to reveal what aspects of the abrupt changes are a result of the hosing itself rather than consequences of the changing AMOC.

2 Experimental setup

2.1 Model description

The simulations shown here use the coupled ocean–atmosphere model CM2Mc, as described in Galbraith et al. (2011). This is a moderately low resolution, but full complexity model, that includes an atmospheric model that is at the high-complexity end of the spectrum applied in previously-published water hosing simulations. In brief, the model includes: a 3° finite volume atmospheric model, similar to that used in the GFDL CM2.1 model (Anderson, 2004); MOM5, a non-Bousinesq ocean model with a fully-nonlinear equation of state, subgridscale parameterizations for mesoscale and submesoscale turbulence, vertical mixing with the KPP scheme as well as due to the interaction of tidal waves with rough topography but otherwise a very low background vertical diffusivity (0.1 cm² s⁻¹), similar to that used in the GFDL ESM2M model (Dunne et al., 2012); a sea-ice module, static land module, and a coupler to exchange fluxes between the components. In addition, the ocean model includes the BLING biogeochemical model as described in Galbraith et al. (2010). This includes limitation of phytoplankton growth by iron, light, temperature and phosphate, as well as a parameterization of ecosystem structure.
2.2 Experimental design

The ten experiments shown here vary in terms of the prescribed atmospheric CO₂, the size of terrestrial ice sheets, and the Earth’s orbital configuration (obliquity and precession). Terrestrial ice sheets were set to either the “preindustrial” extents, or the full LGM reconstruction of the Paleoclimate Model Intercomparison Project 3 (https://pmip3.lsce.ipsl.fr), in which case the ocean bathymetry was also altered to represent lowered sea level, including a closed Bering Strait. Atmospheric CO₂ concentration has a value of either 270 or 180 ppm depending on whether the ice sheets have a preindustrial or glacial extent, respectively. The obliquity was set to either 22.0 or 24.5°, spanning the calculated range of the last 5 My (Laskar et al., 2004). The precessional phase, defined as the angle between the Earth’s position during the Northern Hemisphere autumnal equinox and the perihelion, was set to two opposite positions, corresponding to the positions at which the boreal seasonalities are least and most severe (90 and 270°, respectively). All other boundary conditions of the model were configured as described in Galbraith et al. (2011).

Freshwater forcing (“hosing”) was applied to four simulations run under “Preindustrial” conditions with extrema obliquity and precession values, and another four under “Glacial” conditions also with extrema obliquity and precession. For each of these eight hosing simulations, the model was initialized from a previous preindustrial or LGM state, run for 1000 years, a freshwater hosing was applied for 1000 years, and then it was run for a further 1000 years with the hosing off. Control simulations (without hosing) were also run for the same length of time as the hosing to allow drift correction. During hosing, freshwater was added in the North Atlantic by overriding the land-to-ocean ice calving flux with a preindustrial annual mean plus an additional 0.2 Sv evenly distributed in a rectangle bounded by 40 to 60° N and 60 to 12° W. It should be noted that because the model does not have a rigid lid, this represents a “real” freshwater input to the ocean. The result is a global sea level rise of 15 m over the 1000 year hosing, representing approximately one third the rate of sea level rise estimated for
the maximum rate of sea level rise during the last deglaciation (Melt Water Pulse 1a) (Deschamps et al., 2012). For the top panels of Figs. 2–7, the “weak” AMOC state is defined as years 901–1000 of hosing, and the “strong” AMOC state is the same century taken from the corresponding control simulation to correct for drift, which was quite small.

In addition, we show results from a simulation for which hosing was not applied, but which exhibits spontaneous oscillations in the AMOC reminiscent of D-O events, very similar to those of Peltier and Vettoretti (2014). This “Unhosed” simulation was conducted under glacial CO$_2$ (180 ppm), but with preindustrial ice sheets and bathymetry, low obliquity, and weak boreal seasonality. The final experiment is a modification of the Unhosed simulation, in which a small (0.05 Sv) hosing is applied during a strong-AMOC interval. This final experiment provides an explicit test of the effect of a hosed vs. an unhosed AMOC weakening.

3 Results

3.1 Simulated changes in the North Atlantic

The eight hosed simulations show a number of common features in the North Atlantic, which vary as a function of boundary conditions (Fig. 1). Within the first two centuries of hosing, the Greenland temperature rapidly drops by 8–15 °C, with a stronger response under high obliquity. When the hosing is stopped, the temperatures abruptly increase by 8–20 °C, with a stronger response under high obliquity, and under glacial conditions. The simulated magnitudes of warming and cooling are of the same magnitude as reconstructed for abrupt climate change from Greenland ice cores (Buizert et al., 2014). The AMOC follows a very similar temporal progression in all cases, which is the inverse of the sea ice extent in the North Atlantic. The average North Atlantic sea surface salinity drops by 2–4 PSU in the hosed simulations, consistent with the generally-accepted mechanism of a halocline strengthening due to freshwater forcing, associated with an
expansion of sea ice, being the cause of an AMOC interruption. Thus, as shown by many prior simulations, the hosing response is quite consistent with observations and the idea of iceberg armadas shutting down the AMOC and initiating abrupt climate change (Kageyama et al., 2013; Menviel et al., 2014).

However, Fig. 1 also shows the same metrics of North Atlantic variability for the Unhosed simulation. Under the low CO$_2$, low obliquity and preindustrial ice-sheets of this simulation, the model's coupled ocean–atmosphere system adopts a mode of circulation whereby the AMOC spontaneously oscillates between 12–15 and 6 Sv, on a centennial timescale (Figs. 1 and 2). Although the amplitude of the variations in the Unhosed simulation tends to be smaller, with a weaker AMOC throughout (Fig. 2), and the timescale of fluctuations tends to be shorter than the imposed forcing of the Hosed simulations, the general relationship between the four variables is quite similar (Fig. 1). Thus, the Unhosed oscillations are also consistent with observational evidence that a strengthened halocline, associated with an expansion of sea ice, were coupled to a weakening of the AMOC – even in the absence of external freshwater forcing. A very similar type of unforced AMOC oscillation was observed recently in the Community Earth System (CESM) model, run at higher resolution, by Peltier and Vettoretti (2014).

The Unhosed variability in CM2Mc can be understood as follows. While the AMOC is weak, heat accumulates in the subsurface North Atlantic (500–1500 m) due to the northward transport of heat from the low latitude thermocline, while the strong halocline impedes the downward mixing of cold surface waters. Eventually, the temperature inversion that results from the buildup of heat at depth counteracts the halocline to the point at which the water column can be destabilized. In the model this occurs when a high-salinity anomaly at the sea surface, caused by decadal variability in precipitation and evaporation, as observed in the modern ocean (Glessmer et al., 2014), triggers open-ocean convection. Following destratification, the AMOC rapidly accelerates to its strong state within five years, while the rapid release of heat from the warm subsurface waters melts back the North Atlantic sea-ice area by as much as 80%. The summer sea ice retreat has a large impact on the local radiative balance by changing albedo,
while the winter sea ice retreat allows much more sensible and latent heat release to the atmosphere, so that the Greenland annual average temperature jumps by 10°C within a decade. Over the next decades to centuries, the strong-AMOC state gradually decays back to a state with greater North Atlantic sea-ice, a low AMOC of 6 Sv and cold Greenland temperatures.

The coupled ocean–ice–atmospheric oscillations exhibited by the Unhosed simulation are quite consistent with the mechanism proposed by Kaspi et al. (2004), using a simple box model, as well as by Dokken et al. (2013) based on observations from the Nordic seas. It is remarkable that the magnitude of the salinity change in the unforced simulations is on the order of more than 1 PSU, nearly half as much as for the hosing simulation under the same CO₂ and orbital configuration (Fig. 1), in spite of the absence of any external sources of freshwater. In accordance with the suggestion of Kaspi et al. (2004); Li (2005); Dokken et al. (2013), and the similar simulation of Peltier and Vettoretti (2014), we propose that this type of nonlinear changes in sea ice extent, and their coupling with the oceanic and atmospheric circulation, are at the essence of abrupt climate change.

Next, we show how the global consequences of changes in the AMOC depend on the question of whether or not they are forced by freshwater addition. We do so by exploring a few key atmospheric and oceanic variables that have well-documented responses to abrupt climate change as recorded by paleoclimate records.

### 3.2 Global atmospheric response

Abrupt climate change was initially identified in ice core proxy records of atmospheric temperature, and subsequently extended to temperature variations recorded in multiple proxies from around the world (Clement and Peterson, 2008). Figure 3 summarizes the global surface air temperature change that occurs in response to our simulated AMOC interruptions. Note that, because the AMOC changes vary significantly between the simulations, all anomalies are normalized to an AMOC change of 10 Sv in Figs. 3–7. The first order patterns of surface air temperature change are very similar between
the preindustrial and glacial boundary conditions, as well as in the unhosed simulation. The Northern Hemisphere undergoes general cooling in the extratropics, with greatest cooling in the North Atlantic from Iceland to Iberia, and with cooling extending into the tropics of the eastern Atlantic, North Africa, and southeast Asia. The region of maximum cooling in the Pacific occurs at the Kuroshio-Oyashio confluence, consistent with a southward shift in the front. Meanwhile, the Southern Hemisphere warms everywhere except in the western tropical Pacific and Indian oceans. The maximum warming occurs at high latitudes of the Southern Ocean, and is associated with a sea ice retreat. The general temperature pattern is consistent with the idea of a bipolar seesaw (Broecker, 1998; Stocker and Johnsen, 2003).

One notable difference in the spatial patterns is the temperature change in the NE Pacific, including the western margin of Canada and southern Alaska. The temperature here is quite sensitive to the degree of northward transport of warm ocean water from the subtropical gyre, which is quite variable between simulations. All preindustrial hosed simulations develop a strong Pacific Meridional Overturning Circulation (PMOC) when the AMOC weakens, previously described in models as an Atlantic-Pacific See-saw (Saenko et al., 2004; Okumura et al., 2009; Okazaki et al., 2010; Chikamoto et al., 2012; Hu et al., 2012). The development of a strong PMOC counteracts the hemisphere-wide cooling in the NE Pacific, which can actually cause a warming. This suggests that temperature proxy records from this region would provide strong observational constraints on the degree to which a PMOC developed during abrupt climate changes.

The response to hosing in the North Atlantic is particularly dependant on orbital configurations in glacial simulations because of different initial sea-ice extents which are strongly driven by obliquity (standard deviation in Fig. 3 and Supplement Fig. 1). Otherwise, the response in glacial and preindustrial conditions is fairly consistent between the four extrema orbital configurations.

Changes in precipitation during abrupt climate changes have also been well-documented in speleothems and marine sediment records. As shown in Fig. 4, the
overall patterns of change are similar between Hosed and Unhosed simulations, as they were for the changes in temperature, including reduced precipitation in the North Atlantic, a southward shift of the Intertropical Convergence Zone (ITCZ), and reduced precipitation over south Asia. In fact, many aspects of the precipitation changes vary more as a function of background climate state, including orbital configuration, than they do between Hosed and Unhosed. The two most notable differences between the simulations are the pattern of change surrounding the west Pacific warm pool, an extremely dynamic region with heavy precipitation, and the NE Pacific, where changes in precipitation follow the sea surface temperature through its control on water moisture content and atmospheric circulation. The west Pacific warm pool is an area whose response to hosing is strongly dependant on orbital configuration (standard deviation in Fig. 4), mostly precession (Supplement Fig. 2).

3.3 Global ocean biogeochemistry response

The observed footprint of abrupt climate change also extends to ocean biogeochemistry, with pronounced and well-documented changes in both dissolved oxygen concentrations and export production. Prior work has shown that many aspects of the observed oxygenation changes (Schmittner et al., 2007b) and export production changes (Schmittner, 2005; Mariotti et al., 2012) can be well reproduced by coupled ocean-biogeochemistry models under hosing experiments.

As shown by Fig. 5, our model simulations produce consistent changes in intermediate-depth oxygen during AMOC weakening, hosed or unhosed, that agree well with the bipolar seesaw-mode changes in oxygenation extracted from sediment proxy records of the last deglaciation by Galbraith and Jaccard (2015). All simulations show a decrease of oxygen throughout the full depth of the North Atlantic, due to the lack of ventilated North Atlantic Deep Water, and an increase in the oxygenation of the intermediate-depth North Pacific and Arabian Sea. We note that the North Pacific changes reveal a pronounced shift in the southeastward penetration of North Pacific Intermediate Water, with a hotspot of oxygen change where the edge of the strongly-
ventilated thermocline impinges on the California margin. This hotspot implies that a tendency for a frontal shift to occur in this region makes it particularly sensitive to changes in the AMOC, and explains why the California borderlands region has such rich records of oxygenation changes on millennial timescales.

As shown by Fig. 6, the simulated changes in the North Pacific and Arabian Sea are driven by large changes in ventilation, that overcome basin-wide changes in export production (Fig. 7) acting in the opposite direction (i.e. greater export when intermediate-depth oxygen is higher). The primary discrepancies between the simulations are in the degree of oxygen enrichment in the subarctic Pacific, again related to the development of a PMOC, the changes in ventilation in the Tropical Atlantic thermocline, and in the degree of ventilation changes in the Southern Ocean.

The relative changes in export production, shown in Fig. 7, are locally quite large (in excess of 100 %) but have weaker regional patterns that are less consistent between simulations. The most consistent strong features are a reduction of export in the northern North Atlantic, the western tropical North Atlantic, and the southern margin of the Indo-Pacific subtropical gyre, and an increase in export off of NW Africa, to the west of California, and in the high latitude Southern Ocean. Regions that do not always respond consistently are the subarctic Pacific, which depends significantly on whether or not a PMOC develops, and the Pakistani margin and eastern tropical Pacific, which both show an increase in export under hosing as observed by Kienast et al. (2006) except under glacial hosing, when they do not change. In general, the differences between the various hosed simulations is as large as the difference between hosed and unhosed.

### 3.4 Hosing the unhosed

The analyses above suggest that the large-scale responses of climate and ocean biogeochemistry to an AMOC disruption are quite similar regardless of whether the interruption is forced through hosing or as a spontaneous result of internal model dynamics. In order to further test this apparent insensitivity of results to the trigger of
AMOC weakening, we applied a freshwater forcing during a strong AMOC interval of the Unhosed simulation, forcing the model to return to its weak AMOC state earlier than in the standard Unhosed case. This experiment provides a direct comparison between an unforced and a freshwater forced AMOC reduction. The freshwater forcing weakens the AMOC transport (Fig. 1), and maintains its weakened state over the full 1000 year simulation, whereas the Unhosed simulation returned to the strong-AMOC state after about 800 years. By comparing both the forced and unforced simulations after 800 years stadial, we can estimate the differences caused by the freshwater itself (Fig. 8).

As shown, the changes caused by freshwater itself are quite small compared to the overall changes (compare Fig. 8 with Figs. 3–7 and Supplement Figs. 1–4). Surface air temperatures only differ significantly in the North Atlantic, North Pacific, and high latitude Southern Ocean, while precipitation and oxygen generally show an amplification of the Unhosed trends, and export production shows minimal changes except for the N Atlantic and NE Pacific. Most of these differences can be understood by the fact that the freshwater-forced simulation has a slightly weaker AMOC, and more extensive North Atlantic sea-ice coverage, leading to slightly lower temperatures in the NE Atlantic and a correspondingly greater response in atmospheric circulation that amplifies most features of the weak-AMOC state.

Apart from the small amplification of the general trends, we note one distinct additional feature: the Antarctic temperature response is decreased under hosing, weakening the bipolar seesaw. This feature of the Antarctic response appears to reflect the transport of freshwater from the North Atlantic to the Southern Ocean through relatively shallow pathways in the Atlantic, so that it freshens the Southern Ocean surface. The addition of freshwater strengthens the Southern Ocean halocline, reducing ocean heat release and keeping a larger mantle of sea ice around Antarctica. Thus, the cooler Southern Ocean does not reflect a different response to the AMOC weakening itself, but rather a secondary effect of the freshwater addition through its direct influence on the vertical density structure of the Southern Ocean. We note that Schmittner et al.
(2007a) found an opposite effect of freshwater input on the Southern Ocean, with a relative weakening of the halocline causing a destratification of the Southern Ocean under freshwater forcing. It would thus appear that this aspect of hosing is quite sensitive to model behaviour and experimental design. Given the importance of Southern Ocean convection on controlling atmospheric CO$_2$ (Sarmiento and Toggweiler, 1984; Sigman et al., 2010; Bernardello et al., 2014), this is not a trivial finding.

4 Discussion and conclusions

Our Unhosed model simulation adds to a small, but growing subset of complex 3-dimensional ocean–atmosphere model simulations exhibiting unforced oscillations similar to the abrupt climate changes of Dansgaard-Oeschger cycles. Under a particular set of boundary conditions (low CO$_2$, with preindustrial ice sheets and low obliquity), it can spontaneously switch to a much weaker state, triggered by internal climate variability within the model, acting on a type of “salt oscillator”. This model behavior is essentially identical to that described by Peltier and Vettoretti (2014) and Vettoretti and Peltier (2015).

These spontaneous AMOC variations suggest that when the AMOC is in a weak state due to the background climate state, it is easy for it to switch into an off mode with very little forcing. Under these conditions, large volcanic eruptions (Pausata et al., 2015), ice-sheet topography changes (Zhang et al., 2014) or ocean–atmosphere climate variability controlling precipitation in the North Atlantic, would be sufficient to trigger an abrupt change. However, the ubiquitous weakening of the AMOC that occurs in models under sufficient hosing implies that, even in a strong mode, the AMOC is vulnerable to freshwater forcing, if it is large enough.

As previously suggested, melting of ice shelves could provide an important feedback to an initial AMOC weakening, accelerating ice sheet mass loss due to their buttressing-effect on upstream ice, and adding freshwater that would help stabilize the AMOC in a weak mode (Marshall and Koutnik, 2006; Alvarez-Solas et al., 2010; Marcott et al.,
Such a behaviour is similar to that exhibited here by the Hosed-Unhosed simulation, and may have been necessary to produce a complete “shutdown” of the AMOC, as suggested by Pa/Th measurements at Bermuda Rise (Böhm et al., 2015). This suggests that the question of whether or not a Heinrich event occurred during an AMOC interruption had more to do with the susceptibility of the Laurentide ice sheet to collapse than the nature of the initial AMOC interruption itself. In turn, the degree to which consequent ice sheet melting altered ocean circulation may have depended on where the freshwater was discharged, including how much of the freshwater was input to the ocean as sediment-laden hyperpycnal flows (Tarasov and Peltier, 2005; Roche et al., 2007).

Although they only occur in our model under an unrealistic combination of boundary conditions, a fact which is likely sensitive to the parameterization of deep ocean mixing (Peltier and Vettoretti, 2014), the spontaneous nature of these oscillations allows a powerful comparison to be made with the more typical freshwater-hosed simulations of AMOC weakening. When compared with the Hosed simulations, the general features of the atmospheric and oceanic responses are remarkably robust. These robust features can therefore be taken to reflect consistent dynamical changes related to the AMOC interruption and its coupling with sea ice and atmospheric changes, independent of the ultimate cause of the AMOC interruption. The only substantial difference noted between Hosed and Unhosed simulations lay in the direct impact of freshwater on ocean density structure. The injection of freshwater in the North Atlantic can quickly spread throughout the global ocean, dependent on the distribution of its input and the ocean circulation pattern. In our simulations, the most important consequence of this is the stratification of the Southern Ocean, which is intensified by the freshwater initially injected to the North Atlantic. Aside from this relatively minor contrast, the global features of an AMOC weakening appear to depend just as much on the background climate state as they do on its fundamental cause.
Acknowledgements. We thank the Canadian Foundation for Innovation (CFI) and an allocation to the Scinet supercomputing facility (University of Toronto), made through Compute Canada, for providing the computational resources. We are also grateful towards the open-source communities behind the development of all the python libraries used to create the graphics for this article, particularly Matplotlib (Hunter, 2007) and Cartopy (MetOffice).

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Figure 1. North Atlantic climate metrics for all simulations. The left column shows four simulations under the Glacial boundary conditions (LGM ice sheets and bathymetry with closed Bering Strait; 180 ppm CO$_2$), the central column shows four simulations run under the Preindustrial boundary conditions (Preindustrial ice sheets and bathymetry; 270 ppm CO$_2$), and the right column shows the Unhosed simulation in which spontaneous AMOC variations occur under constant boundary conditions. For the left and center columns, a freshwater hosing of 0.2 Sv was applied between years 1001 and 2000. Orbital configurations are indicated as follows: red = high obliquity; blue = low obliquity; dark = weak boreal seasonality; pale = strong boreal seasonality. The dark grey lines in the right column show the “Hosed-Unhosed” simulation, in which 0.05 Sv of freshwater was added in the North Atlantic during an unforced AMOC “interstadial”. AMOC is defined as the maximum Atlantic streamfunction between 500–5500 m and 30° N : 50° N.
Figure 2. Atlantic meridional overturning streamfunctions in strong and weak states. All plots show 100-year averages. In Glacial and Preindustrial hosing simulations, weak AMOC state is defined by averaging the last century of hosing (model years 1901–2000) and strong AMOC state is defined by averaging the same years from the corresponding control simulation (to correct for potential drift). In the Unhosed simulation, weak AMOC state is defined by averaging the last century of an “unforced” AMOC decrease (model years 4981–5080) and strong AMOC state is defined by averaging a century following an “unforced” AMOC increase (model years 3801–3980). For the top two rows, the streamfunctions are averaged over the corresponding four sets of orbital configurations. The bottom row shows the streamfunctions for the Unhosed simulation. Contours in Sverdrups.
Figure 3. Stadial surface air temperature change normalized to a 10 Sv AMOC decrease. The top two panels show the normalized surface air temperature difference between weak and strong AMOC states (as defined in Fig. 2) under Preindustrial and Glacial boundary conditions, averaged between the four sets of orbital configurations (see Supplement Fig. 1 to view all 8 hosing simulations individually). Light grey contours show the standard deviation between the four sets of orbital configurations at 1 °C intervals. The bottom panel shows the normalized surface air temperature difference between weak and strong AMOC states (as defined in Fig. 2) in the Unhosed simulation. Shading and contours in °C.
Figure 4. Stadial precipitation change normalized to a 10 Sv AMOC decrease. The top two panels show the normalized precipitation difference between weak and strong AMOC states (as defined in Fig. 2) under Preindustrial and Glacial boundary conditions, averaged between the four sets of orbital configurations (see Supplement Fig. 2 to view all 8 hosing simulations individually). Grey contours show the standard deviation between the four sets of orbital configurations at $5 \times 10^{-6}$ kgm$^{-2}$s$^{-1}$ intervals. The bottom panel shows the normalized precipitation difference between weak and strong AMOC states (as defined in Fig. 2) in the Unhosed simulation. Shading and contours in $10^{-6}$ kgm$^{-2}$ s$^{-1}$.
Figure 5. Stadial intermediate-depth (400–1400 m) oxygen concentration change normalized to a 10 Sv AMOC decrease. The top two panels show the normalized oxygen concentration difference between weak and strong AMOC states (as defined in Fig. 2) under Preindustrial and Glacial boundary conditions, averaged between the four sets of orbital configurations (see Supplement Fig. 3 to view all 8 hosing simulations individually). Light grey contours show the standard deviation between the four sets of orbital configurations at $10 \times 10^{-6} \text{ mol kg}^{-1}$ intervals. The bottom panel shows the normalized oxygen concentration difference between weak and strong AMOC states (as defined in Fig. 2) in the Unhosed simulation. Markers represent the 1st principal component of the detrended time series for benthic oxygenation proxies compiled by Galbraith and Jaccard (2015). Shading and contours in $10^{-6} \text{ mol kg}^{-1}$. 
Figure 6. Stadial intermediate-depth (400–1400 m) ideal age change normalized to a 10 Sv AMOC decrease. The top two panels show the normalized age difference between weak and strong AMOC states (as defined in Fig. 2) under Preindustrial and Glacial boundary conditions, averaged between the four sets of orbital configurations (see Supplement Fig. 5 to view all 8 hosing simulations individually). Light grey contours show the standard deviation between the four sets of orbital configurations at 30 year intervals. The bottom panel shows the normalized age difference between weak and strong AMOC states (as defined in Fig. 2) in the Unhosed simulation. Shading and contours in years.
Figure 7. Stadial percent export at 100 m change normalized to a 10 Sv AMOC decrease. The top two panels show the normalized ratio of export between weak and strong AMOC states (as defined in Fig. 2) under Preindustrial and Glacial boundary conditions, averaged between the four sets of orbital configurations (see Supplement Fig. 4 to view all 8 hosing simulations individually). Light grey contours show the standard deviation between the four sets of orbital configurations at 50% intervals. The bottom panel shows the normalized ratio of export between weak and strong AMOC states (as defined in Fig. 2) in the Unhosed simulation. Shading and contours in %.
Figure 8. Difference between the last century (model years 4981–5080) of an 800 year hosed stadial and an 800 year unhosed stadial, under the same boundary conditions. Scale and units for surface air temperature, precipitation, intermediate depth O$_2$ and export are the same as in Figs. 3, 4, 5 and 7, respectively. Note that, unlike Figs. 3, 4, 5 and 7, these values are not normalized to an AMOC decrease.