F. Adolphi & R. Muscheler: “Synchronizing the Greenland ice core and radiocarbon
timescales over the Holocene – Bayesian wiggle-matching of cosmogenic radionuclide
records”. Reply to Reviewers.

First of all we would like to thank both reviewers for their insightful and helpful comments. Below we will address each comment point by point, showing the reviewers comments in black and our response in blue. Changes to the original manuscript are highlighted in bold. At the end of this file you will find a revised version of the manuscript with changes highlighted in yellow.

Reviewer #1:

My overall assessment is that the treatment of the C-14 and Be-10 records is very well done and that the transfer function will be extremely useful to the paleoclimate community. The authors are experts on C-14 and Be-10 and I congratulate them on applying their knowledge to this issue. I recommend publication in Climate of the Past after some minor (but important) revisions and clarifications.

Thank you.

General points

Climate influences on Be-10

The discussion and treatment of climate influences on Be-10 is generally good, however there are a few places where I think it could be improved:

PP2936, L13: It should be noted that most sites receive Be-10 from a combination of wet and dry deposition processes. For example, the detailed treatment of wet and dry deposition processes in the recent paper by Elsässer et al., [2015] suggests.

We agree. In fact, we address this point in line 15-16 (pp2936) by writing “While today wet deposition is dominating over Greenland (Heikkilä et al., 2011) the ratio between wet and dry deposition has likely changed over time (Alley et al., 1995).” Furthermore, by using concentrations and fluxes separately, we account for both endmembers (wet/dry deposition) of these processes. We show that these lead to consistent results (e.g., figure 5). Any mixing between the two modes of deposition would hence, likely not change our results.

We changed PP2936, L13 to:
“In reality, both modes of deposition contribute to the accumulation of 10Be on the ice sheet. Today, wet deposition processes dominate over dry deposition which accounts for about one third or less of the deposited 10Be in Greenland (Heikkinen et al., 2011, Elsaesser et al. 2015). However, this dry/wet deposition ratio has likely been variable over time (Alley et al., 1995).”

PP2936, L18: Please add some more specific language on processes which affect Be-10 concentration in ice e.g. revise to something like “: : :Be-10 transport paths, including stratosphere to troposphere exchange and air-mass precipitation history and can cause climatic imprints...”

We changed PP2936, L18 to: “Secondly, a variety of climatic influences can leave an imprint in ice core 10Be records. Atmospheric circulation changes and air mass precipitation history (i.e., 10Be scavenging by precipitation prior to the arrival of the air mass at the ice core site) may, for example, modulate the transport path and efficiency of 10Be delivery to the ice core site (Heikkinen & Smith 2013, Pedro et al. 2011, 2012). Furthermore, changes in the exchange rates between stratospheric (high 10Be concentrations) and the tropospheric (low 10Be concentrations) air masses can affect the tropospheric 10Be budget (Pedro et al., 2011).”

PP2936, L18: Elsässer et al., 2015 should be added to the list here, they suggest a modest polar bias, similar to the results of Field et al.

We added Elsaesser et al. 2015 to the reference list in PP2936, L 26-27.

PP2942, L22: Note that the Pedro et al., result refers to the coastal East-Greenland ice core site Das2 (not GRIP or GISP).

Yes. However, we feel that PP2942, L17-23 should be seen as a short review of the range of partly disagreeing results that have been published on the issue of a potential polar bias in 10Be records.

PP2945:L5: The statement here assumes that the regression and linear detrending with respect to other proxies in the ice core data does in fact remove all centennial scale climate influences on Be-10. While I agree that the linear detrending is a good step, it is not clear that it would remove all climate influence on Be-10. For example, changes in stratosphere to troposphere exchange are expected to influence Be-10 but not necessarily the other proxies that have been used in the detrending. Please add something to the effect of “a caveat is that
climate influences specific to Be-10 will not be removed by the detrending technique”.

A good point. We changed PP2945, L3-6 to:

“This is in agreement with Adolphi et al. (2014) who showed that centennial GRIP 10Be variations are dominated by solar activity changes and indicate only little climatic influences on 10Be sensitivity to the assumed mode of 10Be deposition even over large deglacial climatic transitions. Other potential climatic influences on 10Be such as changes in the stratosphere-troposphere exchange rates are, however, difficult to assess from climate proxy data and will thus, not be removed by our detrending technique.”

PP2945:L5: “little climatic influence on Be-10 even over large deglacial climatic transitions”. This statement could be misinterpreted. It is well established that glacial to interglacial climate transitions leave a very large imprint on Be-10, mainly due to accumulation rate changes e.g. Finkel and Nishiizumi [1997]. Please revise the wording.

Agreed. We added:

“It should be noted that this statement solely refers to the filtered centennial 10Be variations investigated here.”

PP2945, L3 and Figure 2. “The centennial changes in the GRIP and GISP2 Be-10 versions, however, are highly coherent and indicate a limited climate influence on Be-10 on these timescales”. Clarify if you refer to coherence between the GISP2 and GRIP records or coherence internally within the GISP2 and GRIP records. I would agree that there is good coherence between the curves from the same site but it is not clear that there is good coherence between the records from GRIP and GISP2. A panel in Figure 3 should be added to the main text or at least to the review response showing the Be-10 concentration anomalies at both GRIP and GISP2. The authors do not necessarily need to explain the differences between the records but they should at least be acknowledged given the statement “indicate a limited climate influence on Be-10 on these timescales”. If the authors intend to say that there is good coherence between GRIP and GISP2 records, it could help clarify the section to explain explicitly what is meant by “coherence”, i.e. that the records share the main peaks but not the smaller variations.
We changed figure 3 to the one shown below. To minimize the effect of sampling resolution differences we bandpass-filtered both records [120-500 yrs], i.e., the spectral band that is continuously resolved by both records.

We refer to this figure on P2945, L17:

“Similarly, the $\Delta^{14}$C anomalies modelled from GRIP and GISP2 $^{10}$Be agree within ± 2.5 ‰ (figure 3, c, f).”

To further clarify that P2945, L1-3 refers to each ice core separately we changed it to:

“The centennial changes in the GRIP $^{10}$Be versions, however, are highly coherent and indicate limited climate influence on these timescales and the same holds true for the GISP2 $^{10}$Be versions.”

Figure 3 (revised). Centennial (<500 years) $\Delta^{14}$C variations modelled from GRIP and GISP2 $^{10}$Be data. Panels a and b show the modelled $\Delta^{14}$C variations from $^{10}$Be concentrations (solid black), fluxes (solid grey), “climate corrected” concentrations (dotted black), and “climate corrected” fluxes (dotted grey) for the GRIP (a) and GISP2 (b) $^{10}$Be records. Panels d and e on the right side depict the probability density functions for the maximum $\Delta^{14}$C difference between curves shown in panels a and b, respectively. Panel c shows the mean of all GRIP (black) and GISP2 (grey) $^{10}$Be based $\Delta^{14}$C anomalies shown in panels a and b, respectively. Panel f shows the corresponding probability density function of the maximum $\Delta^{14}$C differences. For this comparison both ice core records have been band-pass filtered [120 –
500 years] to minimize inconsistencies arising from their different sampling resolution. The
correlation between the GRIP and GISP2 records is given in panel c together with its p-value.

**Inferred C-14 production rates and IntCal13–GICC05 transfer function**

I will state up front that I do not have expertise in the Bayesian techniques used here
and hence I cannot critically review that aspect of the methodology. Nevertheless, I
have some questions and comments about the data treatment that I feel are important
to be addressed and which may help make the manuscript accessible to a wider
audience.

- PP2945 L17: Add a panel to Figure 3 showing the comparison between time
  series of the GRIP and GISP based C-14 reconstructions. This would be useful
  for the reader to see directly the coherence between the two, otherwise explain
  why this direct comparison is not needed.

See above. We added a panel to figure 3.

The advantage and influence on uncertainty of going from 50 year \( P_{\text{scaled}} \) to the
annual resolved final age transfer function could be better explained. It would
help to clarify how the interpolation affects the uncertainty if you could plot the \( P_{\text{scaled}}(ts) \)
(i.e on its 50 yr spacing) against the probability distributions for the
final age transfer function. The authors should address if the proposed Monte
Carlo method adds information compared to the simpler approach of interpolating
between the 95% confidence intervals.

We do provide a direct comparison of \( P_{\text{scaled}}(ts) \) to the transfer function in figure 9. The
approach to the transfer function is outlined in section 2.5. The key difference between \( P_{\text{scaled}} \)
and the transfer function is, that \( P_{\text{scaled}} \) contains high-frequency variability in the derived
IntCal13-GICC05 difference which we do not believe we can reliably reconstruct. Since
\( P_{\text{scaled}} \) is estimated on 1,000yr windows every 50 years, neighbouring windows contain about
95% of the same data. Hence, they cannot be used as independent likelihood estimates of the
timescale difference. To account for this oversampling each transfer function samples only
independent \( P_{\text{scaled}} \) estimates (i.e., based on windows that are 1,000 years apart). Each
iteration of the MonteCarlo method starts from a randomly selected \( P_{\text{scaled}} \) estimate and then
uses only every 20\(^{th}\) window (i.e., one \( P_{\text{scaled}} \) estimate per 1,000 years). Hence, every
individual transfer function has a lower sampling resolution than \( P_{\text{scaled}}(ts) \) itself. Through
interpolation of the transfer functions to annual resolution, each likelihood estimate at each
point in time will consist of direct sampling of \( P_{\text{scaled}}(ts) \) at this point in time as well as
interpolated values between neighbouring windows. This can be seen in figure 9 – the
transfer function is much smoother than \( P_{\text{scaled}}(ts) \) but it encompasses the full range of
uncertainty. Consequently, the transfer function has larger uncertainty during periods where
Pscaled implies rapid changes in the timescale difference such as e.g. around 7-8 kaBP, because we cannot reliably resolve the exact rate of change at this point.

Whether we interpolate the transfer functions to annual or 50yr resolution has no effect on the uncertainty, and was merely done for practical reasons, i.e., to be able to provide an annual transfer function to potential users.

Fig 9: The thin black lines are not defined (2-sigma on transfer function?). Also please clarify the difference between the thin black lines and the Pscaled(ts). Why do the Pscaled(ts) and transfer function deviate from each other in places (e.g. 7.5 to 8.0 ka). Is this the result of the 50-yr to annual Monte-Carlo interpolation, which goes back to my question above?

Yes, these are the 2sigma intervals. We changed the figure caption and legend of figure 9 accordingly.

Regarding the disagreement between Pscaled and the transfer function we hope we could explain the effect above.

P2948, eq. 4: The difference between the tree-ring and Be-10 based delta C-14 values is sometimes zero (Fig. 7), and as the IntCal term is always positive, the equality in eq. 4 is not satisfied. The statement may be true if using < sign instead. But the whole section appears a little convoluted. It seems to me that what you do is to adjust the Be-10 scaling factor to minimize the (rms or rmsbinned) difference between the tree-ring and Be-10-based Delta C-14 values.

After obtaining the best value of the scaling factor, you can use the rearranged eq. 4 to estimate (what I would say is the lower bound of) the uncertainty of the Be-10-based Delta C-14 values.

Good point. We agree, our previous formula was not complete and could be misinterpreted. Therefore, we clarified our method by including the rearranged equation 4 into the main text:

\[ \partial(t)_{Be} = \sqrt{\partial(t)^2 - \partial(t)_{IC}^2}; \quad \partial(t) > \partial(t)_{IC} \]

\[ \partial(t)_{Be} = 0; \quad \partial(t) \leq \partial(t)_{IC} \]
We think that this method is more appropriate than simply minimizing the RMSE between 14C and 10Be, since it accounts for the fact that 14C errors are increasing back in time due to the relatively short half-life of 14C.

We cannot evaluate whether this uncertainty estimate is a lower bound of the true uncertainty, but it is the required error to bring 14C and 10Be into statistical agreement, which is crucial for our methodology. However, to clarify that this uncertainty estimate is only valid for the centennial (<500 year) variations in 10Be-based Δ^{14}C we added to P2950, L1-3:

“In conclusion we use a ^{14}C : ^{10}Be ratio of 1.1 : 1 and an uncertainty of 4‰ for the modelled Δ14C record to derive a final IntCal13-GICC05 transfer function in the next section. It should be noted that this uncertainty estimate is only valid for the centennial (<500 year) variations and the period studied here.”

Section 2.2: Please specify whether you stretch or only shift the timescale of the ice core Be-10 data to get an optimal fit with the IntCal C-14 data within each 1000-year window. This may be clear to those familiar with the Bronk Ramsey et al. (2001) paper, but it would be good to make it explicit here.

We changed PP2940, L3 to: “For each window we test for time scale differences (shifts) of ± 150 years without stretching or compression of the time scale within this window.”

Section 2.2: The authors tests the method for robustness in many ways, but the 1000-year width of the correlation window is not tested. That test should be added or at least the authors should discuss why 1000 years is the best choice.

Good point. We did indeed test this effect.

We added on PP2940, L 4: “We tested different window sizes between 500 and 2,000 year length and the corresponding results are consistent within error. The choice of a 1,000 year window represents a trade-off between i) an increasing statistical robustness and hence, smaller uncertainties, and ii) a loss of detail (variability) in the final transfer function (see also section 2.5) with increasing window length.”

Section around P2954, L4. Please note that the main part of the estimated IntCal13-GICC05 difference builds up during the period 8 – 10.3 ka BP, which is the section where the dating is based on GRIP CFA data that have fewer components and lower resolution than the NGRIP dataset employed from 10.3 ka BP downwards [Rasmussen et al., 2006]. The difference curve (Figs. 9-11) levels out in the section between 10.3 ka BP and the onset of the Holocene, corroborating that there are much smaller systematic counting errors in the section based on NGRIP CFA data.

Yes this is correct. We feel that we do point this out in P2954, L 3-6 where we write:
“It can, however, not be assumed that the counting error continues to be systematic beyond this period, since the parameters used for layer identification as well as the sources of uncertainty (e.g. melt layers) differ back in time under changed climatic conditions (Rasmussen et al., 2006).”

Fig 7 and Section 2.2 and 2.5: The fit is very impressive. It would be help the reader to see how your method has reduced uncertainly between the timescales if you could also show some comparisons before synchronisation. I would suggest to show at least one, and preferably 2-3 examples (e.g. best, typical, worst) of 1000-year long sections of wiggle-matched records to allow the reader to evaluate the robustness of the fit.

We added Figure 10 to section 3.4 (see below). We picked the sections between 3,500 – 4,500 BP (a section of relatively low amplitude $\Delta^{14}$C changes, i.e., the variations are close to the estimated 10Be RMSE), between 7,000-8,000 BP (i.e. a section with larger $\Delta^{14}$C variations, but not continuously good agreement between 10Be and 14C), and 10,000-11,000 BP (i.e. a section with large $\Delta^{14}$C variations and a near perfect fit between 10Be and 14C).

We expanded the main text on P2950 (L18 onwards):

“Figure 10 shows three examples of GRIP 10Be based $\Delta^{14}$C anomalies before (grey) and after (black) synchronization to IntCal13 (red). The examples encompass (i) a period of relatively low $\Delta^{14}$C variability ($\pm$5-7‰) but good agreement between GRIP and IntCal13 (figure 10, a), (ii) a period of large $\Delta^{14}$C variability ($\pm$10‰) but less good agreement between GRIP and IntCal13 (figure 10, b), and (iii) a section of large $\Delta^{14}$C ($\pm$10‰) variability and excellent agreement between GRIP and IntCal13 (figure 10, c).

It can be seen, that in all cases the fit between GRIP and IntCal13 is improved when applying the proposed GICC05-IntCal13 transfer function. However, figure 10 (b) also shows, that short periods of disagreement (i.e., around 7,250 – 7,500 years BP) may remain, as they cannot be reliably resolved by our method which matches 1,000-year-long sections. It should, however, be noted that matching these short sections would i) represent a serious violation of the GICC05 counting error which is minimal over these short periods of time ($\pm$6 years at 2$\sigma$ between 7,250 – 7,500 years BP), and ii) not account for the possibility that 10Be and 14C may simply not agree due to the caveats outlined in the introduction. Furthermore, the applied shift of GICC05 in figure 10 (b) leads to an improved agreement between 14C and 10Be after and prior to 7,250 and 7,500, respectively. Hence, we consider it unlikely that for this short period of time the timescale difference deviates significantly from the estimate for the entire window.”
Figure 10. GRIP/GISP2 $^{10}$Be based $\Delta^{14}C$ before (grey) and after (black) synchronization to IntCal13 (red) for the sections a) 3,500-4,500 years BP, b) 7,000-8,000 years BP, c) 10,000-11,000 years BP.

In the final version please specify where the IntCal13–GICC05 transfer function (and relative and absolute uncertainties) will be made available.

We will provide the transfer function as a supplementary file to this paper and on NOAA.

PP2953, L15-20: Worth to specify that the difference is in the direction of systematic over-counting of years.

We added on P2953, L17: “…(i.e., a systematic over-counting of years).

Technical points

- Many of the figures have multiple lines overlain that become hard to distinguish, e.g Fig 5 has 4 lines plus shading. Use of color would probably improve clarity.
- The figures also appear small. I had to zoom in on the screen to see important details. Can you make the figures bigger?

We added color to the lines in Figure 5. Regarding the size of the figures: We will provide them in full A4 size.

In general, ‘both’ is overused. When it is clear that you are talking about two things it is mostly not needed to say both. An example: PP2945, “Both, changes in ocean ventilation [and] air–sea gas-exchange can cause $^{14}C$ anomalies larger than the amplitude of $^{14}C$ anomalies induced by $^{14}$C production rate changes only”. Drop the “both”, it only confuses things here. Also: “One method to compare and synchronize both timescales is the use of cosmogenic radionuclide records”. Here, “both” is misleading unless you are synchronizing (both) time scales to a third one.

Agreed. We reworked the use of “both” carefully.

P2935, L22: “ideal tool” is overstating things given the climate and carbon cycle influences.
We exchanged “ideal” with “powerful”.

P2936, L11: Delete “On the other hand”.

Done.

Section 2.2: Typos: Bronk not Bonk.

Fixed.

Please state what dating of GISP2 was used: obviously it should be GISP2 on GICC05.

On P2938, L7, we added: “We used the GISP2 $^{10}$Be record on the GICC05 timescale (Seierstad et al. 2014).”

Figure 1 Caption: Key data not Key-data. I also noticed some other examples of funny use of hyphens. Please check usage throughout.

Changed figure 1 figure caption and we will check the rest of the manuscript for similar mistakes.

Figure 2 Caption: I can’t make sense of the second last line, please revise.

Ok. Actually this sentence may not be necessary at all, since we discuss the differences between the different $^{10}$Be versions (concentrations, fluxes, climate corrections) later on in Figure 3 and the corresponding text sections. We deleted this sentence from the figure caption.

Figure 5 caption: Description of panel b) appears to be referring to an earlier version.

No, this is indeed the correct caption. To clarify: The patch shows probabilities based on GRIP $^{10}$Be, where gaps in the GRIP record have been filled using GISP2 $^{10}$Be data. Hence, GISP2 can only be used as an independent validation where GRIP and GISP2 have
overlapping $^{10}\text{Be}$ data. The 95% confidence intervals based on GISP2 $^{10}\text{Be}$ during these overlapping sections is plotted as lines in comparison to the patch, which is the same in all panels (GRIP $^{10}\text{Be}$ with GISP2 $^{10}\text{Be}$ filling the gaps). We hope that this is clear from the first 5 lines of the figure caption as well as from P2946, L20-22.

Note some inconsistency in x-axis labels: sometimes yrs BP and sometimes years BP.

**Changed consistently to “years BP”**.

Add space between ka and BP, eg on PP2947.

**Done.**

**Anonymous Referee #2**

The manuscript is well-written and easy to follow. The authors have great attention for detail, which results in realistic uncertainty estimates that reflect all the probable causes of uncertainty.

Thank you.

My only concern is that it’s hard to discern what is really new here. A very similar paper was published last year by the same authors (Muscheler, Adolphi and Knudsen, 2014), MAK14 hereafter. The manuscript under review uses the same $^{10}\text{Be}$ and D14C data as MAK14, and while the mathematical details differ, the approach is conceptually identical (i.e. converting ice-core $^{10}\text{Be}$ to 14C using a carboncycle model, and then wiggle-matching it to IntCal D14C). Unsurprisingly, the transfer function the authors derive is essentially identical to the one derived by MAK14 – only smoother due to the choice of a 1000 year window length. The main improvement is a reduction in the uncertainty estimates, suggesting that MAK14 were too conservative in estimating their error.

The aim of this manuscript is twofold.

1. We present a novel approach to synchronize radionuclide records from different archives that is applicable to a multitude of records. We anticipate that the methodology itself will be applied in further studies, and hence, we describe in detail its underlying assumptions and
caveats. In addition to MAK14 we also test our results using different 10Be records (GRIP and GISP2) to illustrate the robustness of our results and methodology. Furthermore, we assess uncertainties arising from the geochemistry of 10Be and 14C explicitly which is a significant improvement compared to MAK14 which also has general implications for the interpretation of these records (i.e. for solar activity reconstructions). Hence, we think we can provide a significant conceptual advance on how to link cosmogenic radionuclide records and lie out a framework that will be of great value for future studies.

2. We provide a transfer function between GICC05 and IntCal13 for the Holocene. Compared to MAK14 this function is less (not more, see figure 7 in MAK14) smooth than the proposed transfer function of MAK14 which is based on a 2,000 year window. More importantly, we do not only provide smaller, but also more robust uncertainty estimates on the transfer function (which we have to acknowledge were not satisfactorily defined in MAK14). This is a significant improvement as reliable uncertainty estimates are crucial if the transfer function shall be applied to determine robust leads and lags in the climate system between ice core and radiocarbon dated paleoclimate records.

We clarified the value of this publication with respect to MAK14 by modifying the abstract to:

„Compared to earlier work, we employ a novel statistical approach which leads to strongly reduced and yet, more robust, uncertainty estimates. Furthermore, we demonstrate that the inferred timescale differences are robust independent of (i) the applied ice core $^{10}$Be records, (ii) assumptions of the mode of $^{10}$Be deposition, as well as (iii) carbon cycle effects on $^{14}$C, and (iv) in agreement with independent estimates of the timescale differences.“

In addition, section 1.1 (aim of this study) refers to the advantages of the proposed method and tests as compared to MAK14.

The work is very thorough, and I have only a few minor comments that should be addressed in a revised manuscript. I leave out the first two digits (“29”) in all listed page numbers.

- I think section 2.2 (statistical method) would fit more logically between the current sections 2.4 and 2.5. When reading the section on the statistical method, the reader has no idea what is meant by “$^{10}$Be-based D$^{14}$C anomalies” (P38, last line). This becomes clear after reading section 2.4. An alternative solution would be to add an introductory paragraph to section 2.2 in which the conceptual framework is laid out, so the reader understands that $^{10}$Be is converted to $^{14}$C using a carbon cycle model, and then filtered to isolate the centennial component.

This is a good point. We added an introductory paragraph to section 2.2:
“In the following section we will describe the statistics involved in the $^{14}$C/$^{10}$Be comparison. To be able to compare both radionuclides quantitatively, we converted the ice core $^{10}$Be records into $\Delta^{14}$C variations using a box-diffusion carbon cycle model (Siegenthaler et al., 1980; Muscheler et al., 2004b). The details of this conversion and its uncertainties are addressed in more detail in section 2.4. In the following we will refer to the modelled $\Delta^{14}$C variations as ,,$^{10}$Be-based $\Delta^{14}$C anomalies“.”

- Due to their proximity, the GISP2 and GRIP sites should experience identical atmospheric $^{10}$Be loading; yet GISP2 receives slightly more accumulation than GRIP (about 5%). Could this help in partitioning out wet and dry $^{10}$Be deposition? The lower $^{10}$Be concentrations at GISP2 (by 0.12 atoms/g), as well as the higher $^{14}$C/$^{10}$Be scaling factors (Figs 6 and 7) are both consistent with a fraction of dry deposition. I fear that the accumulation difference may be too small to do this reliably, though.

This is an interesting thought. However, to assess this reliably a detailed evaluation of the $^{10}$Be concentration / ice accumulation rate relationship would be required. For the purpose of clarity of the manuscript we would prefer not to discuss this issue extensively. We show, that GRIP and GISP2 $^{10}$Be records yield similar synchronization results to IntCal13 (figure 5, b). Furthermore, $^{10}$Be concentrations and fluxes give consistent synchronization estimates (figure 5, a) indicating that the assumed mode of $^{10}$Be deposition is of minor importance for the results of this study. Regarding the slightly different $^{10}$Be scaling factors of GRIP and GISP2 records (figure 7), we think that this difference should not be over-interpreted. The GISP2 $^{10}$Be record has a lower sampling resolution and a slightly higher scaling factor may just result from this difference in smoothing. Last but not least, due to slight differences in the $^{10}$Be sample preparation of both ice cores (see Finkel and Nishiizumi 1997, JGR) we cannot exclude that the small difference in the $^{10}$Be concentration reflects an interlaboratory difference.

- Section 2.1: please indicate the data resolution for the D14C data also.

We added in P2938, L18: “…and presented in IntCal13 in 5-year resolution while the underlying data has typically a resolution of 10 years for most of the Holocene”

- P37 L1-L4: I think the reality is more fluid than portrayed. I suspect that in practice the $^{10}$Be-$^{14}$C synchronization is dominated by a few prominent events, and therefore somewhat “discrete”. Likewise, continuous (rather than discrete) CH4 synchronization has also been achieved between ice cores (Mitchell et al., Science 342 964-966, 2013).
We think that a strength of the methodology applied here lies in its lower sensitivity to single events than the methodology of MAK14. The methodology of MAK14 is based on correlation analysis, which relies on the covariance of two records. Hence, this type of analysis can be dominated by single events of large amplitudes. In comparison, the methodology applied in the manuscript presented here does compare 14C/10Be data pairs in a reduced Chi2-like fashion. I.e., the values are compared directly irrespective of their covariance. This allows for example also to exploit the information “a 500 year long section of zero D14C anomaly” – while the covariance of such a section would be zero. Therefore, common variability as well as common non-variability go with similar weight into the comparison. Obviously, D14C anomalies are needed to achieve a synchronization and (in the example above) constrain the length of a section with small D14C anomalies. And we agree, that a slight dominance of these larger D14C anomalies can be expected, since these anomalies will exceed the uncertainty of the 14C/10Be records, maximizing equation 1. However, as outlined above we do think, that also relatively “flat” sections of 14C/10Be contribute significantly to the synchronization i.e. we evaluate all available information.

To tone down our statement in P2937, L 1-4, and also with respect to the findings by Mitchell et al. (2013), we changed P2937, L 1-4 to:

“... has the advantage that it can provide near-continuous estimates of the time scale differences [...] or changes in atmospheric trace gases during Dansgaard-Oeschger events.”

- P40, L25: “....and may thus diminish the climate influences in the 10Be record”. Could it also increase the climate influences in the 10Be record, if the observed correlations are spurious?

Yes, in theory this could happen. However, the removed correlations were indeed significant. In addition, we hope we explain sufficiently well on P2940 (L26) – P2941 (L5) that these corrections are not meant to be interpreted as improved “climate free” versions of the 10Be records, but as sensitivity tests to our methodology.

- Section 2.4.1: What is the motivation to only investigate the sensitivity of the model to the oceanic carbon exchange? While the ocean is of course the largest carbon reservoir, the terrestrial carbon fluxes are actually larger than the oceanic ones. A recent paper also suggested that changes in terrestrial carbon reservoirs are more important during Holocene (Bauska et al. Nat Geo 8, 383-387 2015)

We agree that terrestrial carbon fluxes are a major component in the global carbon cycle. However, due to the short turnover rate of the terrestrial biosphere, the biologically stored carbon has essentially the same 14C signature as the atmosphere. Hence, changes in the biosphere – atmosphere CO2 exchange, do not exert a strong control on atmospheric 14C concentrations (i.e. a large flux of terrestrial carbon won’t change the $^{14}\text{C}/^{12}\text{C}$ ratio in the atmosphere significantly).
Section 2.6: I think it’s important somewhere to point out that you’re comparing the 14C anomalies, rather than 14C itself. These anomalies are not really well defined; from section 3.1 I assume you’re using the centennial (<500 yr) variations. Please describe how your filter the records to separate the <500 and >500 yr variations.

We explain in section 2.2 (P2938, L23-26) why we use D14C anomalies rather than absolute values or 14C-ages. We consider the use of a 500 year high-pass filter a result rather than part of the method as it results from the climate and carbon cycle related 10Be uncertainties (section 3.1) which are timescale dependent and minimize for these short wavelengths.

We specified more clearly on P2946, L13:

“In the following we will compare the centennial (i.e., <500 years separated by an FFT-based High-Pass filter) D14C anomalies as reconstructed from tree-rings (IntCal13) and ice cores (GRIP/GISP2 10Be-based) with respect to their timescale differences. The choice of a 500 year high-pass filter results from the climate and carbon cycle related uncertainties shown in section 3.1 which increase on longer timescales.”

-P45, L16-17: How much is the uncertainty of 3 ‰ relative to the standard deviation of the data itself? In other words, what is the signal to noise ratio?

The standard deviation of the GRIP 10Be based D14C record is 4.5 per mille, yielding a signal to noise ratio of ca. 1.5. We’d like to point out that the uncertainty estimate is very conservative since we treat the climate influence on 10Be as a systematic uncertainty (i.e., Figure 3 shows the difference between the maximum and minimum D14C value at each point in time, as opposed to the standard deviation of all 10Be versions).

-P48, L21: “this would imply a strong polar bias”. Please elaborate, this is not automatically clear.

We introduce and discuss the issue of a polar bias in the introduction and in section 2.4.2. To make it more clear we changed L21 on P2948 to: “Assuming that the centennial 10Be and 14C production rate changes are mainly modulated through solar activity this low scaling factor would point to a strong polar bias of the GRIP GISP2 10Be records (see sections 1 and 2.4.2).

-The generated transfer function should be provided as a text / excel file in the supplement.

Yes, it will be.
- Typos / Language:

P35 L15 and throughout: acronym should be capitalized, so GCR instead of gcr.

Done.

P36, L4: 14C / 12C *ratio*

Added.

P38, L23: Please define what is meant by “D14C anomalies”. I don’t think this is done anywhere in the manuscript.

We define D14C on P2938, L16 and think that the term “anomalies” is a general term that describes deviations from a mean. As outlined above, we think that the use of centennial anomalies (i.e., deviations from a 500 year low-pass filtered D14C record) is a result of section 3.1, which may partly differ if the method was applied to a different 10Be record, with different climate influences. We do outline on P2938, L23-26 why absolute D14C values cannot be used and hope that our definition of anomalies becomes clear throughout the manuscript and with the additions on P2946, L13 (see above).

To increase the clarity of the manuscript we added on P2938, L 26:

“Given the results shown in section 3.1 we employ centennial (<500 year high-pass filter) Δ\textsuperscript{14}C anomalies of the tree-ring and the 10Be-based Δ\textsuperscript{14}C records for this comparison as shown in figure 3.”

P39, L4 and L8: Bronk Ramsey (“r” omitted)

Fixed.

Throughout there are long sentences that would benefit from inclusion of a comma to clarify sentence structure. Some examples:

P35 L23: After production, ... P36 L11: On the other hand, ... P36 L28: ... synchronization tools, ... P40 L19: ... to the ice sheet, ... P41 L18: ... these effects, ... P45 L26: ... as before, ...

Reworked.
Synchronizing the Greenland ice core and radiocarbon timescales over the Holocene - Bayesian wiggle-matching of cosmogenic radionuclide records

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Abstract

Investigations of past climate dynamics rely on accurate and precise chronologies of the employed climate reconstructions. The radiocarbon dating calibration curve (IntCal13) and the Greenland ice core chronology (GICC05) represent two of the most widely used chronological frameworks in paleoclimatology of the past ~50,000 years. However, comparisons of climate records anchored on these chronologies are hampered by the precision and accuracy of both timescales. Here we use common variations in the production rates of $^{14}$C and $^{10}$Be recorded in tree-rings and ice cores, respectively, to assess the differences between both timescales during the Holocene. Compared to earlier work, we employ a novel statistical approach which leads to strongly reduced and yet, more robust, uncertainty estimates. Furthermore, we demonstrate that the inferred timescale differences are robust independent of (i) the applied ice core $^{10}$Be records, (ii) assumptions of the mode of $^{10}$Be deposition, as well as (iii) carbon cycle effects on $^{14}$C, and (iv) in agreement with independent estimates of the timescale differences. Our results imply that the GICC05 counting error is likely underestimated during the most recent 2,000 years leading to a dating bias that propagates throughout large parts of the Holocene. Nevertheless, our analysis indicates that the GICC05 counting error is generally a robust uncertainty measurement but care has to be taken when treating it as a nearly Gaussian error distribution. The proposed IntCal13-GICC05 transfer function facilitates the comparison of ice core and radiocarbon dated paleoclimate records at high chronological precision.
1 Introduction

Paleoclimatology can provide significant insights into natural climate changes and thus, improve our understanding of the climate system. Besides the reconstruction of past climate itself, a precise chronology of each paleoclimate record is crucial to reliably assess the dynamics of the inferred changes. Furthermore, consistent chronologies across multiple paleoclimate records are required to assess the spatiotemporal evolution of climatic events and thus, to test for potential leads and lags within the climate system and ultimately improve the understanding of the underlying processes of past climate change. Two independent key timescales in paleoclimatology of the past 50,000 years are the radiocarbon- (IntCal13, Reimer et al., 2013) and the Greenland ice core timescale (GICC05, Andersen et al., 2006; Rasmussen et al., 2006; Seierstad et al., 2014; Svensson et al., 2008; Vinther et al., 2006). To be able to infer leads and lags between paleoclimatic changes anchored on these chronologies at high precision, it is crucial to test the consistency between the timescales and establish climate-independent isochrones and thus, reduce the influence of their absolute dating uncertainties (e.g., Lane et al., 2013). One method to compare and synchronize different timescales is the use of cosmogenic radionuclide records, such as $^{10}$Be and $^{14}$C (Muscheler et al., 2014a; Muscheler et al., 2014b; Muscheler et al., 2008; Southon, 2002).

Cosmogenic radionuclides such as $^{10}$Be and $^{14}$C are produced in the atmosphere through a nuclear cascade mainly triggered by incoming galactic cosmic rays (GCR, Lal and Peters, 1967). The flux of GCR reaching the atmosphere is in turn modulated by the strength of the helio- and geo- magnetic fields resulting in varying production rates of $^{10}$Be and $^{14}$C (Masarik and Beer, 2009, 1999; Kovaltsov et al., 2012; Kovaltsov and Usoskin, 2010). Thus, increased (decreased) intensity of the solar- and/or geomagnetic field will result in decreased (increased) cosmogenic radionuclide production rates. Therefore, $^{14}$C and $^{10}$Be production rates co-vary globally due to external processes, making them a powerful synchronization tool.

After production, $^{14}$C oxidizes to $^{14}$CO$_2$ that enters the global carbon cycle and gets stored in various environmental archives such as tree rings, sediments, and speleothems. $^{10}$Be attaches to aerosols which are deposited within 1-2 years (Raisbeck et al., 1981) by wet and dry deposition processes and is stored in sediments including polar ice sheets. These ‘system effects’ (i.e., non-production influences on $^{10}$Be and $^{14}$C records such as the mixing, transport, and deposition of $^{14}$C and $^{10}$Be) can challenge an unequivocal reconstruction of cosmogenic
radionuclide production rates from paleoarchives and thus, synchronization efforts based on cosmogenic radionuclides.

Due to the large actively exchanging carbon reservoirs, changes in the atmospheric $^{14}$C/$^{12}$C ratio are attenuated and delayed compared to the corresponding $^{14}$C production rate variations (Oeschger et al., 1975). In comparison, $^{10}$Be is a more direct recorder of production rate changes. Thus, when comparing $^{14}$C and $^{10}$Be records directly, this difference in geochemistry has to be taken into account by using carbon cycle models (Muscheler et al., 2004b). However, to be fully realistic, these corrections would require prior knowledge on the variable state of the carbon cycle, which is often difficult to quantify (Köhler et al., 2006).

$^{10}$Be records (for example from ice cores) can be affected by non-production related processes as well. Firstly, it depends on the assumed mode of deposition (wet vs. dry) whether the $^{10}$Be concentration (all wet deposition) or the $^{10}$Be flux (all dry deposition) is the better measure of atmospheric $^{10}$Be concentration changes (Alley et al., 1995; Delaygue and Bard, 2010). In reality, both modes of deposition contribute to the accumulation of $^{10}$Be on the ice sheet. Today, wet deposition processes dominate over dry deposition which accounts for about one third or less of the deposited $^{10}$Be in Greenland (Heikkilä et al., 2011; Elsässer et al., 2015). However, this dry/wet deposition ratio has likely been variable over time (Alley et al., 1995). Secondly, a variety of climatic influences can leave an imprint in ice core $^{10}$Be records. Atmospheric circulation changes and air mass precipitation history (i.e., $^{10}$Be scavenging by precipitation prior to the arrival of the air mass at the ice core site) may, for example, modulate the transport path and efficiency of $^{10}$Be delivery to the ice core site (Heikkilä and Smith, 2013; Pedro et al., 2012; Pedro et al., 2011b). Furthermore, changes in the exchange rates between stratospheric (high $^{10}$Be concentrations) and the tropospheric (low $^{10}$Be concentrations) air masses can affect the tropospheric $^{10}$Be budget (Pedro et al., 2011a). Thirdly, contrary to $^{14}$C, $^{10}$Be might not be hemispherically well mixed owing to its short atmospheric residence time. This has led to the proposition of a so-called “polar bias” in ice core $^{10}$Be records, stating that if polar $^{10}$Be records were dominated by $^{10}$Be produced at high latitudes, the anisotropy of the geomagnetic shielding would lead to an enhanced solar- and an attenuated geomagnetic modulation signal in polar $^{10}$Be records. There is contradicting evidence from data and modelling studies to whether this is the case (Field et al., 2006; Bard et al., 1997; Pedro et al., 2012; Muscheler and Heikkilä, 2011; Heikkilä et al., 2009; Elsässer et al., 2015).
In summary, to be able to use $^{10}\text{Be}$ and $^{14}\text{C}$ as synchronization tools, ‘system effects’ on each radionuclide have to be assessed and corrected for. If successful, this method has the advantage that it can provide near-continuous estimates of time scale differences as opposed to discrete tie-points obtained from tephrochronology (Abbott and Davies, 2012; Lane et al., 2013) or changes in atmospheric trace gases during Dansgaard-Oeschger events (Blunier et al., 1998; Buizert et al., 2015).

1.1 Aim of this study

Recently, Muscheler et al. (2014a) assessed the differences of the radiocarbon and ice core time scales for the past 14,000 years by comparing GRIP $^{10}\text{Be}$ (Yiou et al., 1997; Muscheler et al., 2004b; Vonmoos et al., 2006) and IntCal13 $^{14}\text{C}$ data (Reimer et al., 2013). Here, we revisit this approach using a different statistical framework (Bronk Ramsey et al., 2001) that is computationally less expensive and provides improved error estimates for the inferred timescale differences as compared to the method used in Muscheler et al. (2014a). Furthermore, we test the robustness of the obtained results with respect to the use of different ice core $^{10}\text{Be}$ records as well as potential ‘system effects’ on the radionuclide records. We focus our analysis on the period where dendrochronologically dated high quality $^{14}\text{C}$ measurements on tree rings are available. While this is theoretically the case back to 12,560 calBP (calibrated before present, AD1950, Friedrich et al., 2004), the accuracy of the oldest part of tree-ring chronology has recently been questioned (Hogg et al., 2013) causing a gap in the $^{14}\text{C}$ records underlying IntCal13 around 12,000 calBP (Reimer et al., 2013). Hence, we limit our analysis to the Holocene where dendrochronological and $^{14}\text{C}$-data replication is high and most robust (Reimer et al., 2013; Friedrich et al., 2004).

2 Methods

2.1 Data

The key data used in this paper is shown in figure 1. The GRIP $^{10}\text{Be}$ record (Vonmoos et al., 2006; Muscheler et al., 2004b; You et al., 1997) covers almost the entire Holocene with a gap between 9,400 and 10,800 years BP (Before Present 1950 AD) and no data for sections younger than 300 years BP. We use the data as presented in Vonmoos et al. (2006) that
includes a 61-point binomial filter (roughly corresponding to a 20 year low-pass filter or a
decal sampling resolution) minimizing weather related noise in the $^{10}$Be data. The GISP2
$^{10}$Be record (Finkel and Nishiizumi, 1997) has a gap between 7980 and 9400 years BP and no
data for sections younger than 3270 years BP. We used the GISP2 $^{10}$Be record on the
GICC05 timescale (Seierstad et al., 2014). Its temporal resolution varies between 20 to 60
years with an average of one sample every 35 years. Hence, no smoothing filter was applied.
The GISP2 $^{10}$Be concentrations have been normalized to the same standard used for the GRIP
$^{10}$Be measurements (NIST SRM 4325, see Yiou et al., 1997; Muscheler et al., 2004b). The
resulting GRIP and GISP2 $^{10}$Be records differ by on average 0.12 $10^4$ atoms/g of ice. To avoid
inhomogeneities when splicing the records together, we adjusted the GISP2 $^{10}$Be data
accordingly by adding 0.12 $10^4$ atoms/g to the GISP2 $^{10}$Be record (see figure 1). We note that
reconciling the $^{10}$Be records through normalization instead of addition does not affect the
results shown here. The lower panel in figure 1 shows atmospheric $\Delta^{14}$C (that is $^{14}$C/$^{12}$C after
correction for fractionation and decay relative to a standard) as reconstructed from
dendrochronologically dated tree rings (Friedrich et al., 2004) and presented in IntCal13 in 5-
year resolution while the underlying data has typically a resolution of 10 years for most of the
Holocene (Reimer et al., 2013).

### 2.2 Statistical method

In the following section we will describe the statistics used for the $^{14}$C/$^{10}$Be comparison. To
be able to compare both radionuclides quantitatively, we converted the ice core $^{10}$Be records
into $\Delta^{14}$C variations using a box-diffusion carbon cycle model (Siegenthaler et al.,
1980; Muscheler et al., 2004b). The details of this conversion and its uncertainties are
addressed in more detail in section 2.4. In the following we will refer to these modelled $\Delta^{14}$C
variations as “$^{10}$Be-based $\Delta^{14}$C anomalies”.

We employ a statistical approach that is commonly used in the ‘wiggle-match dating’ of $^{14}$C
records that have an initial relative chronology, i.e. the age differences between neighbouring
samples are known, such as tree-rings (Bronk Ramsey et al., 2001). Contrary to classical $^{14}$C-
age calibration we use $\Delta^{14}$C anomalies, since $^{10}$Be cannot provide information on absolute
$\Delta^{14}$C (and hence, $^{14}$C-ages) which depends on $^{14}$C production rates and the state of the carbon
cycle long before the investigated period. Given the results shown in section 3.1 we employ
centennial (<500 year FFT high-pass filter) $\Delta^{14}$C anomalies of the tree-ring and the $^{10}$Be-
based $\Delta^{14}C$ records for this comparison as shown in figure 3. The mathematical formulation remains however, unchanged. The calibration record, IntCal13 (Reimer et al., 2013), describes $\Delta^{14}C$ anomalies for each point in time, $R(t)$, with an associated uncertainty, $\delta R(t)$. This can be compared to $^{10}Be$-based $\Delta^{14}C$ anomalies ($R_i;\pi$) for which we know the absolute age differences ($\Delta t_i$) between each sample from ice core layer counting. We can estimate the probability ($P_i$) for different assumed time scale differences between the records ($t_s$) for each sample by using equation 8 in Bronk Ramsey et al. (2001):

$$P_i(t_s + \Delta t_i) \propto \frac{\exp\left(-\frac{(R_i - R(t_s + \Delta t_i))^2}{2(\delta R_i^2 + \delta R^2(t_s + \Delta t_i))}\right)}{\sqrt{\delta R_i^2 + \delta R^2(t_s + \Delta t_i)}}$$  \hspace{1cm} (1)

Using Bayes’ theorem to combine the probabilities for each individual measurement we can obtain an overall probability ($P_s$) for each time scale difference between GICC05 and IntCal13 (equation 9 in Bronk Ramsey et al., 2001):

$$P_s(t_s) \propto \prod_{i=1}^{n} P_i(t_s + \Delta t_i)$$  \hspace{1cm} (2)

To allow a continuous comparison, all records have been interpolated to annual resolution. However, since the ice core sampling resolution is in reality lower we do not obtain truly independent probability distributions for each sample. Consequently, we correct for the reduced degrees of freedom by scaling $P_s$ as:

$$P_{s\text{scaled}}(t_s) = P_s(t_s)^{1/r}$$  \hspace{1cm} (3)

where $r$ is the original sample spacing (years/sample) of the ice core $^{10}Be$ records. This scaling effectively widens the obtained probability distribution and thus, increases the derived uncertainties. For the filtered GRIP $^{10}Be$ record, we assume a decadal resolution.

This ‘wiggle-matching’ is done for predefined windows of IntCal13 and GRIP and hence, yields a probability distribution ($P_{s\text{scaled}}(t_s)$) for their time scale difference for each window. We apply this method to 1,000 year windows of $^{14}C/^{10}Be$ data and investigate one window every 50 years back in time. For each window we test for time scale differences (shifts) of $\pm$ 150 years without stretching or compression of the timescale within this window. Hence, in analogy to $^{14}C$-wiggle-match dating, each window could be seen as a single 1,000 year long “tree” that is being calibrated. We tested different window sizes between 500 and 2,000 year length and the corresponding results are consistent within error. The choice of a 1,000 year window represents a trade-off between (i) an increasing statistical robustness and hence,
smaller uncertainties, and (ii) a loss of detail (variability) in the final transfer function (see also section 2.5) with increasing window length.

It can be seen from equation 1, that contrary to the correlation analysis employed by Muscheler et al. (2014a) this method favours $^{10}$Be/$^{14}$C linkages with a direct 1:1 relationship between IntCal13 and $^{10}$Be-based Δ$^{14}$C records. Hence, the $^{14}$C:$^{10}$Be production rate ratio has to be assessed. Furthermore, the uncertainty for the $^{10}$Be-based records and the $^{10}$Be:$^{14}$C conversion is quantitatively included in the calculation and hence, needs to be estimated. In the following sections we will outline how these factors can be initially assessed.

2.3 Assessment of uncertainties due to climatic influences on $^{10}$Be

As outlined in the introduction, ice core $^{10}$Be records can be affected by various climatic influences that can ‘contaminate’ the production signal. To account for these effects, we use four different versions of the GRIP and GISP2 $^{10}$Be records throughout the manuscript. We use $^{10}$Be concentrations and fluxes ($^{10}$Be concentration multiplied by snow accumulation and ice density) as endmembers of the assumed mode of $^{10}$Be deposition (wet vs. dry, respectively) on the ice sheet. To address the role of climate influences on $^{10}$Be mixing and transport to the ice sheet, we additionally generated “climate corrected” versions of the concentrations and fluxes. For this purpose, we performed multiple linear regression analysis between $^{10}$Be and climate proxy time series from the GRIP and GISP2 ice cores. Using ice accumulation rates (Seierstad et al., 2014), δ$^{18}$O (Johnsen et al., 1995; Stuiver et al., 1997), and ion data (Mayewski et al., 1997) as predictors, we linearly detrended the $^{10}$Be concentrations and fluxes. This procedure removes covariance between $^{10}$Be and climate proxy data and may thus, diminish the climate influences in the $^{10}$Be record. It should be noted, that this is a ‘blind’ empirical approach that does not aim for a process based understanding of the climate influences on $^{10}$Be. This method would, for example, confound solar ($^{10}$Be) variations that had an influence on climate as climate influences on $^{10}$Be (Adolphi et al., 2014). Hence, these ‘climate corrected’ versions should rather be seen as sensitivity tests for our analysis than as improved estimates of past $^{10}$Be production rates per se. In summary, we use four (concentrations, fluxes, and “climate corrected” versions thereof) different versions of the GRIP and GISP2 $^{10}$Be data. Each version represents a plausible endmember of the $^{10}$Be production rate history, depending on the assumed mode of
deposition and climatic impacts on $^{10}$Be and can thus, be used to assess the sensitivity of our analysis to these processes.

2.4 Assessment of uncertainties due to $^{10}$Be – $^{14}$C conversion

2.4.1 Carbon cycle modelling

To be able to compare $^{10}$Be to $^{14}$C records, we converted the $^{10}$Be records into $\Delta^{14}$C anomalies using a box-diffusion carbon cycle model (Oeschger et al., 1975; Siegenthaler et al., 1980). The model was run under pre-industrial conditions and has been shown to yield consistent results with more complex carbon cycle models for our purposes (Muscheler et al., 2007). As outlined in the introduction, the unknown state and dynamics of the carbon cycle introduce uncertainty to the comparison of $^{10}$Be and $^{14}$C. To test for the sensitivity to these effects, we conducted four experiments (table 1). Each experiment was forced with an idealized 200 year $^{14}$C production rate cycle of ± 20% approximately corresponding to a solar de Vries cycle. For two of the experiments we perturbed the state of the carbon cycle by increasing (S1) or decreasing (S2) the air-sea gas exchange constant by 50% mimicking changes in wind speed and/or sea ice extent. In the scenarios S3 and S4 the ocean diffusivity parameter (ocean ventilation) was increased and decreased by 50%, respectively. Each experiment was spun up for 50,000 years under preindustrial conditions until all $^{14}$C reservoirs were in steady state. Subsequently the investigated parameter was changed linearly from its preindustrial to its perturbed value within 50 years (transition 1). The perturbed state was then maintained for 25,000 years to reach equilibrium again (steady state) before linearly changing the perturbed parameter back to preindustrial values within 50 years (transition 2). We use these different sensitivity experiments to obtain an uncertainty estimate of the modelled ($^{10}$Be-based) $\Delta^{14}$C records due to carbon cycle effects.

2.4.2 $^{10}$Be/$^{14}$C production rate ratio

To compare tree ring and ice core radionuclide records we used the normalized $^{10}$Be records as $^{14}$C production rate input for the carbon cycle model. This yields a $^{10}$Be-based $\Delta^{14}$C anomaly record that can be directly compared to the tree-ring data. Hence, we have to assume a ratio between the production rates of $^{14}$C and $^{10}$Be. This ratio depends on the radionuclide
production cross sections and the energy spectrum of the incoming GCR. Model estimates ofelative $^{14}$C: $^{10}$Be production rate increases for a change in the solar modulation parameter
from 700 to 0 MeV at modern geomagnetic field strength differ between 1.34 (Masarik and
Beer, 2009) and 1.04 (Kovaltsov et al., 2012; Kovaltsov and Usoskin, 2010). Similarly, the
predicted $^{14}$C: $^{10}$Be production rate ratios for changes in the geomagnetic field strength are
model dependent for unresolved reasons (Cauquoin, 2014).

Furthermore, the $^{14}$C: $^{10}$Be production rate ratio depends on the presence of a potential ‘polar
bias’ (see introduction). If a ‘polar bias’ was present (Bard et al., 1997; Field et al., 2006) the
ratio between $^{14}$C and ice core $^{10}$Be variations could be biased towards lower values. (Bard et
al. (1997) report a value of 0.65 for the South Pole $^{10}$Be record). For Greenland, however,
high resolution $^{10}$Be records do not support such a strong polar bias but would instead be
consistent with a well mixed atmosphere (Pedro et al., 2012; Muscheler and Heikkinä, 2011).
Simply comparing the standard deviations of centennial variations of IntCal13 and $^{10}$Be-
based $\Delta^{14}$C anomalies leads to ratios between 0.95 and 1.05 ($\sigma^{14}$C_{IntCal}/$\sigma^{14}$C_{10Be}) depending on
which ice core (GRIP/GISP2) and which version of the $^{10}$Be records (concentration, flux,
climate corrections) is used. Thus, we start with a $^{14}$C: $^{10}$Be production rate ratio of 1:1 and
test the sensitivity of our results to this assumption by repeating the calculations outlined in
section 2.2 using $^{14}$C: $^{10}$Be ratios of 1.5:1 and 0.5:1.

2.5 Timescale transfer function

The methodology outlined in section 2.2 yields a probability estimate of the IntCal13-
GICC05 timescale difference every 50 years. These probability distributions are however not
fully independent since neighbouring 1,000 year windows overlap and are, hence, largely
based on the same data. To create a timescale transfer function we employed a Monte-Carlo
procedure that creates 20,000 possible transfer functions based on independent, i.e. non-
overlapping, windows. Each iteration, i) randomly selects one of the youngest (most recent)
20 windows and ii) randomly samples from the probability distribution $P_{\text{scaled}}(t_x)$ of this
window as well as the older non-overlapping windows (i.e. one window every 1,000 years so
that the selected windows are fully independent with respect to the data points they contain).
The resulting transfer functions are then interpolated to annual resolution and converted into
probability distributions for the timescale difference at each point in time. For each transfer
function we assume that both timescales are correct at 0 BP (i.e. AD 1950).
2.6 Iterative structure of the synchronization method

The separate aspects of our synchronization method outlined above are applied in an iterative manner to obtain robust and self consistent error estimates for our results. The different steps involved are carried out in the following order:

i. We create four versions of both ice core $^{10}$Be records as endmembers of plausible $^{10}$Be production rate histories (see section 2.3).

ii. We convert these $^{10}$Be records into $\Delta^{14}$C using a box-diffusion carbon cycle model (section 2.4.1) assuming a $^{14}$C:$^{10}$Be production rate ratio of 1 (see section 2.4.2).

iii. The difference between the different $^{10}$Be-based $\Delta^{14}$C records, and results from the carbon cycle sensitivity experiments (see section 2.4.1) serve as initial uncertainty estimates for the $^{10}$Be-based $\Delta^{14}$C records.

iv. We then compare the tree ring and $^{10}$Be-based $\Delta^{14}$C records with respect to their timescale differences using the statistics outlined in section 2.2. We test for the robustness of these results by using all four different $^{10}$Be versions of GRIP and GISP2 separately as well as $^{10}$Be-$^{14}$C conversion factors of 0.5 and 1.5 (see section 2.4.2).

v. Calculating an initial timescale transfer function (see section 2.5) we then synchronize IntCal13 and GICC05. This enables us to directly compare tree ring and $^{10}$Be-based $\Delta^{14}$C records and estimate the optimal $^{14}$C:$^{10}$Be production rate ratio, as well as uncertainties for the $^{10}$Be-based $\Delta^{14}$C record.

vi. Based on these posterior estimates of the $^{14}$C:$^{10}$Be ratio and the uncertainty of the $^{10}$Be records, we repeat the calculations outlined in sections 2.2 and 2.5 yielding our final estimates of the IntCal13-GICC05 timescale differences over the Holocene.

3 Results

3.1 Climate and Carbon cycle related uncertainties in the GRIP and GISP2 $^{10}$Be records

Figure 2 displays the different $^{10}$Be production rate scenarios from GRIP (top two panels) and GISP2 (lower two panels) $^{10}$Be concentrations (Conc), fluxes (Flux) and their climate
corrected versions (\text{Conc}_{\text{clim}} \text{ and } \text{Flux}_{\text{clim}}, \text{ respectively}). Dividing the \(^{10}\text{Be}\) records into a centennial (<500 years) and millennial (>500 years) variations indicates that the different \(^{10}\text{Be}\) versions mainly differ in the low frequency range. These millennial differences can systematically affect the modelling of \(\Delta^{14}\text{C}\) since the carbon cycle acts as an integrator over \(^{14}\text{C}\) production rate variations. The centennial changes in the GRIP \(^{10}\text{Be}\) versions, however, are highly coherent and indicate a limited climate influence on \(^{10}\text{Be}\) on these timescales and the same holds true for the GISP2 \(^{10}\text{Be}\) versions. This is in agreement with Adolphi et al. (2014) who showed that centennial GRIP \(^{10}\text{Be}\) variations are dominated by solar activity changes and indicate only little sensitivity to the assumed mode of \(^{10}\text{Be}\) deposition even over large deglacial climatic transitions. It should be noted that this statement solely refers to the filtered centennial \(^{10}\text{Be}\) variations investigated here. Other potential climatic influences on \(^{10}\text{Be}\) such as changes in the stratosphere-troposphere exchange rates are, however, difficult to assess from climate proxy data and will thus, not be removed by our detrending technique. Thus, in the following we will focus on centennial (<500 years) changes in \(^{10}\text{Be}\) and \(^{14}\text{C}\) production rates to avoid systematic errors originating from uncertainties in the millennial \(^{10}\text{Be}\) production rate history.

The left hand panels in figure 3 show the corresponding modelled \(\Delta^{14}\text{C}\) anomalies from the centennial \(^{10}\text{Be}\) variations indicated in figure 2 assuming a \(^{14}\text{C}:^{10}\text{Be}\) production rate ratio of 1:1. As expected, similar to the \(^{10}\text{Be}\) records these variations are highly coherent. The right panels in figure 3 display histograms of the maximal \(\Delta^{14}\text{C}\) difference between the different production rate histories (i.e. the absolute \(\Delta^{14}\text{C}\) difference between the highest and the lowest modelled \(\Delta^{14}\text{C}\) version at each point in time). It can be seen that the different \(^{10}\text{Be}\) versions translate into a modelled \(\Delta^{14}\text{C}\) uncertainty of about \(\pm 3\ \%\) (1\(\sigma\)) for GRIP (figure 3 a, d) and GISP2 (figure 3 b, e). Similarly, the \(\Delta^{14}\text{C}\) anomalies modelled from GRIP and GISP2 \(^{10}\text{Be}\) agree within \(\pm 2.5\ \%\) (1\(\sigma\), figure 3 c, f).

As outlined in the introduction, the state and the dynamics of the carbon cycle impose an uncertainty on the \(^{10}\text{Be}-^{14}\text{C}\) comparison that is difficult to quantify from the data itself (Köhler et al., 2006; Muscheler et al., 2004b). Figure 4 shows the results from the performed carbon cycle sensitivity experiments (see section 2.4.1, table 1). It can be seen that the millennial \(\Delta^{14}\text{C}\) variations are substantially altered by carbon cycle perturbations (figure 4 b). Changes in ocean ventilation (experiments S3 and S4) and well as air-sea gas exchange (experiments S1 and S2) can cause \(\Delta^{14}\text{C}\) anomalies larger than the amplitude of \(\Delta^{14}\text{C}\) anomalies induced by \(^{14}\text{C}\) production rate changes only (control). However, as before, the
centennial $\Delta^{14}$C variations are considerably less affected by these perturbations (figure 4 c).

The increase (decrease) of air-sea gas exchange or ocean ventilation does lead to a decrease (increase) in the amplitude of the modelled centennial $\Delta^{14}$C variations. However, these changes in amplitude are largely limited to about $\pm 3 \%_e$ (figure 4, panel d) except for about 200-300 years around the timing of the carbon cycle perturbation itself (figure 4, transitions 1 and 2). Importantly, the phase of the centennial $\Delta^{14}$C variations is not affected by the imposed carbon cycle changes. Since the applied carbon cycle changes in our sensitivity experiments are likely unrealistically large for Holocene conditions (Köhler et al., 2006; Roth and Joos, 2013), we conservatively assume a 1σ uncertainty of $\pm 3 \%_e$ (see figure 4, panel d, ‘steady state’) for the modelled $\Delta^{14}$C records due to carbon cycle effects.

Adding the uncertainties due to climate impacts on $^{10}$Be ($\pm 3 \%_e$) and the carbon cycle ($\pm 3 \%_e$) in quadrature we thus, obtain an initial uncertainty estimate of about $\pm 4.5 \%_e$ for the modelled $\Delta^{14}$C records.

### 3.2 Sensitivity of the synchronization method to uncertainties in the $^{10}$Be-$^{14}$C conversion

In the following we will compare the centennial $\Delta^{14}$C (i.e., $<500$ years, separated by an FFT-based high-pass filter) anomalies reconstructed from tree rings (IntCal13) and ice cores (GRIP/GISP2 $^{10}$Be-based) with respect to their timescale differences. The choice of a 500 year high-pass filter results from the climate and carbon cycle related uncertainties shown in section 3.1 which increase on longer timescales. We use the statistical framework outlined in section 2.2 and assign an initial uncertainty of $\pm 4.5 \%_e$ to the $^{10}$Be-based $\Delta^{14}$C records. The uncertainties for the tree-ring based $\Delta^{14}$C anomalies are taken from IntCal13 (Reimer et al., 2013). For this purpose we spliced the GISP2 $^{10}$Be versions into the corresponding GRIP $^{10}$Be versions to fill the gap in the GRIP record between 9,400 and 10,800 years BP and create a continuous record for the entire Holocene. Hence, in the following “GRIP” refers to this combination of GRIP and GISP2 data, while results for the GISP2 data are only shown for periods where they have not been used to fill the gap in the GRIP record.

Figure 5 displays the obtained probability distributions $P_{\text{scaled}}(t_s)$ for each sliding window, centred on its mean age. The results are shown for all four GRIP $^{10}$Be versions (panel a), in comparison to results based on GISP2 data only (panel b), as well as for different assumed
$^{14}$C:$^{10}$Be production rate ratios (panel c). The different GRIP $^{10}$Be versions yield consistent estimates of the IntCal13-GICC05 timescale differences throughout the Holocene. The only marked difference occurs around the 8.2 ka BP event (Blockley et al., 2012). During this period the $^{10}$Be flux indicates a more rapid increase in the IntCal13-GICC05 timescale difference as compared to all other $^{10}$Be versions. As noted by Muscheler et al. (2004a) the accumulation rate anomaly associated to the climate oscillation around 8,200 years ago appears to lead to an ‘over correction’ of the $^{10}$Be deposition during flux calculation. This leads to a worse agreement between $^{14}$C and $^{10}$Be fluxes as compared to $^{14}$C and $^{10}$Be concentrations (see figure 3 in Muscheler et al., 2004a). This is corroborated by the fact that results based on the “climate corrected” $^{10}$Be flux follow the probability estimates of $^{10}$Be concentrations (figure 5a).

Comparing GRIP based results to GISP2 based estimates indicates consistent estimates of the timescale differences. The larger uncertainties of the GISP2 based results are due to the lower sampling resolution of the GISP2 $^{10}$Be record (see equation 3).

Figure 5c shows the sensitivity of our results to the assumed $^{14}$C:$^{10}$Be production rate ratio. It can be seen that the inferred timescale differences are relatively insensitive to the assumed $^{14}$C:$^{10}$Be ratio. However, the derived uncertainty of $P_{\text{scaled}}(t_x)$ does increase with lower $^{14}$C:$^{10}$Be ratios. This can easily be understood by imagining a scaling of zero for the $^{10}$Be-based record which would result in an infinitely wide probability distribution.

In summary, our method of estimating the IntCal13-GICC05 timescale difference is i) largely robust for all versions of the GRIP $^{10}$Be record, ii) consistent for GRIP and GISP2 $^{10}$Be data, and iii) independent of the assumed $^{14}$C:$^{10}$Be production rate ratio. However, this analysis also shows that it is important to compare $^{10}$Be concentrations and fluxes to identify potential caveats as seen around the 8.2 ka BP event. Furthermore, while the estimate of the most likely timescale difference (i.e. the location of the maximum of $P_{\text{scaled}}(t_x)$) may not be affected by the assumed $^{14}$C:$^{10}$Be ratio, the uncertainty of this estimate is. Hence, in the following section we will derive a posterior estimate of the $^{14}$C:$^{10}$Be ratio, as well as a refined uncertainty estimate of the $^{10}$Be-based $\Delta^{14}$C records.
3.3 Posterior estimate of the $^{14}$C:$^{10}$Be production rate ratios and uncertainties

As shown in the previous section, our estimates of the most likely timescale difference between IntCal13 and GICC05 are largely independent of which $^{10}$Be record (GRIP/GISP2) and which version thereof (concentration, flux, climate corrections) is used, as well as which $^{14}$C:$^{10}$Be ratio is assumed. Hence, we calculated an initial GICC05-IntCal13 transfer function (section 2.5) and synchronized the tree ring based and $^{10}$Be-based $^{14}$C record. This enables us to compare the records with respect to the most likely $^{14}$C:$^{10}$Be ratio. In addition, we can derive a posterior estimate of the modelled $^{10}$Be-based $^{14}$C uncertainty.

After synchronization we can compare tree ring and $^{10}$Be-based $^{14}$C sample pairs assuming different $^{10}$Be scaling factors (i.e. $^{14}$C:$^{10}$Be ratios) between zero and two. The difference between tree ring and $^{10}$Be-based $^{14}$C sample pairs ($\delta(t)$) is a function of the uncertainty of IntCal13 ($\delta_{IC}(t)$) and the uncertainty of the $^{10}$Be-based records ($\delta_{Be}(t)$) in the form that:

$$\delta(t) = \sqrt{\delta(t)^2_{IC} + \delta(t)^2_{Be}}$$

Hence, we can rearrange equation 4 and use the quoted uncertainties of IntCal13 to derive $\delta(t)_{Be}$:

$$\delta(t)_{Be} = \sqrt{\delta(t)^2 - \delta(t)^2_{IC}}; \quad \delta(t) > \delta(t)_{IC}$$

$$\delta(t)_{Be} = 0; \quad \delta(t) \leq \delta(t)_{IC}$$

These uncertainties can be summarized to the root mean square error (RMSE$_{10Be}$). This way we can obtain the optimal $^{10}$Be scaling factor (where the RMSE$_{10Be}$ minimizes) and the associated uncertainty of the $^{10}$Be-based $^{14}$C records (the minimum of the RMSE$_{10Be}$). Figure 6 displays the results of this analysis indicating an optimal $^{10}$Be scaling factor of around 0.7. Assuming that the centennial $^{10}$Be and $^{14}$C production rate changes are mainly modulated through solar activity this low scaling factor would point to a strong polar bias of the GRIP GISP2 $^{10}$Be records (see sections 1 and 2.4.2). However, when investigating the $^{14}$C time series it becomes apparent, that this low scaling leads to an underestimation of the amplitude of virtually all grand solar maxima and minima (i.e. large $^{14}$C anomalies) in the $^{10}$Be-based $^{14}$C record (figure 7, top). This bias is induced by the fact, that the $^{14}$C anomalies are normally distributed around 0 ‰ leading to a majority of the $^{14}$C values lying close to zero dominating the RMSE$_{10Be}$. Hence, for these values a low scaling of the $^{10}$Be-based $^{14}$C records will simply act to reduce noise from the record and thus, reduce the RMSE$_{10Be}$. 


To avoid this bias, we performed a binned regression analysis. We divided the tree ring and $^{10}\text{Be}$-based $\Delta^{14}\text{C}$ sample pairs into bins of 2.5 ‰ (defined based on the tree ring $\Delta^{14}\text{C}$ anomalies) and calculated the RMSE$_{^{10}\text{Be}}$ for each bin (RMSE$_{^{10}\text{Be}_\text{bin}}$). These uncertainties for each bin can then be summarized to an overall RMSE$_{^{10}\text{Be}}$ as:

$$RMSE_{^{10}\text{Be}} = \sqrt{\text{RMSE}_{^{10}\text{Be}_\text{bin}}^2}$$ (6)

This binning leads to an equal weighting of small and large $\Delta^{14}\text{C}$ anomalies in the comparison of the $\Delta^{14}\text{C}$ records. It can be seen that this method indicates a larger $^{14}\text{C:$^{10}\text{Be}$ ratio of about 1.1 (figure 8) and avoids the systematic underestimation of large amplitude $\Delta^{14}\text{C}$ anomalies (figure 7, bottom). Depending on the production rate model used, this scaling indicates a weak (Masarik and Beer, 2009, 1999) or no (Kovaltsov et al., 2012; Kovaltsov and Usoskin, 2010) polar bias in the Greenland $^{10}\text{Be}$ records. In addition, it can be seen that the minimum of the RMSE$_{^{10}\text{Be}}$ becomes larger than without binning, indicating an uncertainty of about 4 ‰ for the $^{10}\text{Be}$-based $\Delta^{14}\text{C}$ records. This is due to the above described effect, that the noise is not artificially suppressed and can be seen by comparing the decadal scale peaks in the top and bottom panels of figure 7. The larger $^{10}\text{Be}$ scaling factor makes the $^{10}\text{Be}$ record appear noisier. However, firstly, this noise may represent remaining influences of ‘system effects’ on ice core $^{10}\text{Be}$ records and hence, represent an uncertainty that has to be taken into account. Secondly, it should be kept in mind that IntCal13 is a stack of multiple $^{14}\text{C}$ datasets which will inevitably result in smoothing. This smoothing may also reduce the amplitude of ‘real’ $\Delta^{14}\text{C}$ variations instead of merely reducing noise, since the differences between the underlying raw data sets of IntCal13 are potentially in part systematic (Stuiver et al., 1998; Adolphi et al., 2013).

In conclusion we use a $^{14}\text{C:$^{10}\text{Be}$ ratio of 1.1:1 and an uncertainty of 4 ‰ for the modelled $\Delta^{14}\text{C}$ record to derive a final IntCal13-GICC05 transfer function in the next section. It should be noted that this uncertainty estimate is only valid for the centennial (<500 year) variations studied here.

### 3.4 IntCal13-GICC05 transfer function

Using the estimated $^{14}\text{C:$^{10}\text{Be}$ ratio of 1.1 and a $^{10}\text{Be}$-based $\Delta^{14}\text{C}$ error of ±4 ‰ ($\pm$1σ) (see previous section) we recalculated the ‘wiggle-match’ probability distributions ($P_{\text{scaled}}(t_s)$,
equation 3) for the IntCal13-GICC05 timescale difference (figure 9, grey shading). For these
calculations we used the mean of all GRIP-$^{10}\text{Be}$-based $\Delta^{14}\text{C}$ versions (concentration, flux,
climate corrections) and filled the gap between 9,400 and 10,800 yrBP using the GISP2 data.
Based on these probability distributions we modelled the IntCal13-GICC05 transfer function
as described in section 2.5. The resulting transfer function (figure 9 solid lines) averages out
some short-term fluctuations in the timescale difference compared to the initial ‘wiggle-
match’ probability distributions. As described in section 2.5 this is due to the used window
length of 1,000 years to determine $P_{sscaled}(t_s)$ at each point in time, preventing an
independent assessment of faster changes in the timescale difference. Nevertheless, the
estimated uncertainties of the timescale transfer function (thin black lines in figure 9) encompass the uncertainties of the ‘wiggle-match’ probability distribution at each point in
time.

Figure 10 shows three examples of GRIP $^{10}\text{Be}$ based $\Delta^{14}\text{C}$ anomalies before (grey) and after
(black) synchronization to IntCal13 (red). The examples encompass (i) a period of relatively
low $\Delta^{14}\text{C}$ variability ($\pm 5-7\%$o) but good agreement between GRIP and IntCal13 (figure 10, a),
(ii) a period of large $\Delta^{14}\text{C}$ variability ($\pm 10\%$o) but less good agreement between GRIP and
IntCal13 (figure 10, b), and (iii) a section of large $\Delta^{14}\text{C}$ ($\pm 10\%$o) variability and excellent
agreement between GRIP and IntCal13 (figure 10, c). It can be seen, that in all cases the fit between GRIP and IntCal13 is improved when applying the proposed GICC05-IntCal13
transfer function. However, figure 10 (b) also shows, that short periods of disagreement (i.e.,
around 7,250 – 7,500 years BP) may remain, as they cannot be reliably resolved by our
method which matches 1,000 year-long sections. It should, however, be noted that matching
these short sections would (i) represent a serious violation of the GICC05 counting error
which is minimal over these short periods of time ($\pm 6$ years at $2\sigma$ between 7,250 – 7,500
years BP), and (ii) not account for the possibility that $^{10}\text{Be}$ and $^{14}\text{C}$ may simply not agree due
to the caveats outlined in the introduction. Furthermore, the applied shift of GICC05 in figure
10 (b) leads to an improved agreement between $^{14}\text{C}$ and $^{10}\text{Be}$ after and prior to 7,250 and
7,500, respectively. Hence, we consider it unlikely that for this short period of time the
timescale difference deviates significantly from the estimate for the entire window.
4 Discussion

Figure 11 shows the obtained estimate of the IntCal13-GICC05 timescale difference in comparison to the results obtained by using the method of Muscheler et al. (2014a, re-run with a 1,000 year window length) and age markers that have been independently anchored on both timescales.

Our results are fully consistent with the results obtained by Muscheler et al. (2014a). While this is expected to some extent, as our study and the work by Muscheler et al. (2014a) are based on the same data, it shows that the statistical approach used here leads to similar results as the Monte-Carlo lag-correlation analysis but is computationally much less expensive. Furthermore, as shown in figure 5, we obtain similar results when using the GISP2 $^{10}\text{Be}$ instead of the GRIP $^{10}\text{Be}$ record lending additional support to the robustness of our results.

The additional modelling of the transfer function employed here (sections 2.5 and 3.4) leads to a smoother development of the timescale difference which is more realistically reflecting limitations of the method imposed by the window size of the $^{14}\text{C}$-$^{10}\text{Be}$ comparison. The difference between the timescale transfer functions around 8,200 years BP is induced by the fact that Muscheler et al. (2014a) based their calculations on $^{10}\text{Be}$ fluxes which are influenced by accumulation rate changes around this time as discussed in section 3.2 and in Muscheler et al. (2004a).

The largest difference between the results presented here and by those of Muscheler et al. (2014a) is seen in the derived error estimates. We obtain strongly reduced uncertainties for the estimated timescale differences. This is likely due to the fact, that Muscheler at al. (2014a) used a comparably ad-hoc and highly conservative method to derive their uncertainties. By taking the distribution of the mean $r^2$-values of all iterations Muscheler et al. (2014a) do not include the results of the Monte-Carlo analysis of the “Best Fits” in their error estimate. Thus, $^{14}\text{C}$-$^{10}\text{Be}$ matches that may not be the most likely solution in any of the iterations become included in the uncertainty envelope. In comparison, the statistics employed here allow a direct analytical assessment of the synchronization uncertainties. Hence, while our uncertainty estimates are significantly smaller, we consider them more robust. Theoretically, systematic errors from undetected biases in the $^{10}\text{Be}$ record could lead to erroneous results. However, the results shown in section 3.2 demonstrate the consistency of GRIP and GISP2$^{10}\text{Be}$-based calculations as well as for different climate corrections and do, thus, not indicate such biases (see figure 5). In conclusion, while largely consistent, we
regard the method employed here a significant improvement to the approach by Muscheler et al. (2014a).

Comparing our results to independent estimates of IntCal13-GICC05 timescale differences further supports our analyses (figure 11, symbols). Two major solar proton events (“775 and 994 AD events”) leaving well defined spikes in the $^{14}$C content of dendrochronologically dated trees (Miyake et al., 2013; Miyake et al., 2012; Gütter et al., 2015) as well as in Greenland ice core $^{10}$Be records (Mekhaldi et al., 2015; Sigl et al., 2015) indicate an IntCal13-GICC05 timescale difference of $-7 \pm 2$ (2σ) years for both events (Sigl et al., 2015). Consistent with these findings, we obtain IntCal13-GICC05 differences of $-4 \pm 4$ and $-6 \pm 5$ years (2σ) for the 994 and 775 AD event, respectively. It should be noted that these annual radionuclide excursions are not present in the data used here, which is of lower resolution, and are hence, independent estimates of the timescale difference.

Based on tephra findings in the GRIP ice core (Barbante et al., 2013) the historically dated AD 79 eruption of Vesuvius has been used as a reference point in the GICC05 chronology (Vinther et al., 2006). However, our results indicate a timescale offset of $-11 \pm 6$ (2σ) years at AD 79 (1871 years BP, see figure 11). Assuming that the tree-ring chronologies are correct at this time, this would imply an age of AD 90 $\pm 6$ for the GRIP tephra layer – incompatible with an attribution to the age of the Vesuvius eruption within 2σ. This result is in agreement with the analysis by Sigl et al. (2015) who recently counted annual layers in the NEEM and NEEM-2011-S1 ice cores and dated this marker horizon to AD 87 and 89, respectively.

The age of the Minoan eruption of Santorini has long been debated and the presence of an unequivocally attributable signal in the ice core records has been questioned (Pearce et al., 2004; Hammer et al., 1987; Hammer et al., 2003; Friedrich et al., 2006). The GICC05 age of 3591 $\pm 5$ BP of an identified tephra horizon is incompatible with the radiocarbon based age of 3563 $\pm 14$ calBP of the Santorini eruption ($\Delta = -28 \pm 15$ yrs). Our results indicate a chronology difference of $-20 \pm 5$ years around this time, reconciling the two aforementioned ages (see figure 11, open diamond). Hence, at least from a chronological point of view, it cannot be ruled out that the ice core tephra may be ascribable to the Santorini eruption (Muscheler, 2009).

Volcanic glass shards from the Saksunarvatn ash have been found in the GRIP ice core (Grönvold et al., 1995), as well as in multiple marine, lacustrine and terrestrial sites, of which the Lake Kråkenes record provides the highest resolution radiocarbon based age for the
The dating difference of -86 ± 35 years between the radiocarbon based age by Lohne et al. (10,210 ± 35 calBP, ±1σ) and the GICC05 age (10,296 BP, Abbott and Davies, 2012) of the Saksunarvatn ash is consistent with our estimated timescale difference of -66 ± 10 years during this time interval.

In summary, our results are consistent within uncertainties with all independent age markers that link the GICC05 and IntCal13 timescales over the Holocene.

Figure 12 displays the inferred IntCal13-GICC05 timescale differences in comparison to the GICC05 maximum counting error (Rasmussen et al., 2006; Vinther et al., 2006). Assuming that the tree-ring chronologies underlying IntCal13 are accurate throughout the Holocene our results imply an underestimation of the absolute dating uncertainty of GICC05 for large parts of the Holocene. Furthermore, it can be seen that the counting error appears to be systematic, in that most uncertain years (counted as 0.5 ± 0.5 years, Rasmussen et al., 2006) have indeed not been true calendar years during the Holocene (i.e., a systematic over-counting of years). Nevertheless, when comparing the rate of change of the inferred IntCal13-GICC05 timescale difference to the rate of change of the maximum counting error (i.e. the relative maximum counting error) it can be seen that – even though systematic – the identification of uncertain years in the ice core records is accurate. Except for the most recent 2,000 years where (potentially erroneous) fix-points like the Vesuvius eruption are used to constrain GICC05 the relative layer counting uncertainty appears to be an accurate uncertainty estimate. This can be seen in figure 12 (lower panel) which indicates that the rate of change of the GICC05 maximum counting error is consistent within error with the rate of change of the IntCal13-GICC05 timescale difference prior to 2,000 years BP. This is important to note as it generally supports the GICC05 layer counting methodology and uncertainty which forms the basis of GICC05 back to 60,000 years BP (Svenssson et al., 2008), even though the systematic nature of the derived timescale differences challenges the use of the maximum counting error as a nearly Gaussian distributed 2σ uncertainty during the Holocene (Andersen et al., 2006). It can, however, not be assumed that the counting error continues to be systematic beyond this period, since the parameters used for layer identification as well as the sources of uncertainty (e.g. melt layers) differ back in time under changed climatic conditions (Rasmussen et al., 2006).

Alternatively, uncertainties in the dendrochronologies underlying IntCal13 could contribute to the growing discrepancy between IntCal13 and GICC05 over the Holocene. This appears, however, unlikely since the tree-ring chronologies have been cross-dated back to 7,272 calBP.
to the Irish Oak Chronology (Pilcher et al., 1984) and back to 9,741 calBP using independently constructed German Oak Chronologies (Friedrich et al., 2004; Spurk et al., 2002). Furthermore, the gradual development of the timescale difference appears consistent with a counting uncertainty, while a dendrochronological mismatch could be expected to cause sudden ‘jumps’ in the timescale difference. However, consistently missing tree rings in both German oak chronologies for the period older than 7,272 calBP could theoretically contribute to the growing timescale difference.

5 Conclusions

We employed a novel approach to infer timescale differences between two of the most widely used chronologies in Holocene paleoclimatology, the radiocarbon (IntCal13, Reimer et al., 2013) and Greenland ice core (GICC05, Svensson et al., 2008) timescales. Our results are largely consistent with the results of Muscheler et al. (2014a) but yield significantly smaller and more robust uncertainty estimates. The inferred timescale differences are consistent with independent tie-points obtained from volcanic tephras and solar proton events. However, in agreement with Sigl et al. (2015) our analyses indicate that the attribution of an ice core tephra to the AD 79 eruption of Vesuvius (Barbante et al., 2013) may be erroneous which leads to a propagating ice core dating bias that affects large parts of the Holocene. Nevertheless, the identification of uncertain years in the ice core during the Holocene is otherwise generally accurate as expressed in the relative counting error (figure 12 lower panel). This is important to note as it, in principle, supports the layer counting method and uncertainty estimates also beyond the period investigated here. Furthermore, it should be noted that these conclusions are based on the assumption that the tree-ring time scale is accurate.

Independent of the accuracy of either of the two chronologies we provided a high precision transfer function between the radiocarbon and Greenland ice core timescales. This allows radiocarbon dated and ice core paleoclimate records to be compared at high chronological precision which will improve studies of leads and lags within the climate system throughout the Holocene (Bronk Ramsey et al., 2014). Furthermore, the methodology outlined here can be applied to link high resolution $^{14}$C records such as floating tree-ring chronologies to ice core time scales and thus, aid in testing and improving the glacial radiocarbon dating calibration curve.
The proposed GICC05-IntCal13 transfer function shown in figure 9, 11 and 12 is available as a supplementary file to this paper and on NOAA.

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Table 1. Performed carbon cycle sensitivity experiments. All percentage values refer to the control simulation under pre-industrial conditions.

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Figure 1: Top: GRIP (grey, Vonmoos et al., 2006) and GISP2 (black, Finkel and Nishiizumi, 1997) Holocene $^{10}$Be concentrations. The GRIP $^{10}$Be record is smoothed by a 61-pt binomial filter (see Vonmoos et al., 2006). The GISP2 $^{10}$Be record has been shifted by $+0.12\times10^4$ atoms/g to correct for a difference in the mean of the GRIP and GISP2 $^{10}$Be records. Bottom: Atmospheric $\Delta^{14}$C as reconstructed from tree rings (Reimer et al., 2013 and references therein).
Figure 2: Comparison of $^{10}$Be fluxes and concentrations over the Holocene. Solid black and grey curves denote $^{10}$Be concentrations and fluxes, respectively. Dotted lines refer to the “climate corrected” (see text) versions of concentrations and fluxes with similar colour coding as solid lines. The top two panels show GRIP $^{10}$Be for variations on time scales longer (top) than 500 years, and for wavelengths between 100-500 years (below). The 100 year cut-off has been applied for clarity of the figure. The bottom two panels show GISP2 $^{10}$Be for the same wavelengths as for GRIP.
Figure 3. Centennial (<500 years) $\Delta^{14}$C variations modelled from GRIP and GISP2 $^{10}$Be data. Panels a and b show the modelled $\Delta^{14}$C variations from $^{10}$Be concentrations (solid black), fluxes (solid grey), “climate corrected” concentrations (dotted black), and “climate corrected” fluxes (dotted grey) for the GRIP (a) and GISP2 (b) $^{10}$Be records. Panels d and e on the right side depict the probability density functions for the maximum $\Delta^{14}$C difference between curves shown in panels a and b, respectively. Panel c shows the mean of all GRIP (black) and GISP2 (grey) $^{10}$Be based $\Delta^{14}$C anomalies shown in panels a and b, respectively. Panel f shows the corresponding probability density function of their maximum $\Delta^{14}$C differences. For this comparison both ice core records have been band-pass filtered [120 – 500 years] to minimize inconsistencies arising from their different sampling resolution. The correlation between the GRIP and GISP2 records is given in panel c together with its p-value.
Figure 4. Carbon cycle sensitivity experiments. a) Normalized $^{14}$C production rate input to the model. b) Modelled $\Delta^{14}$C anomaly. c) Centennial (<500 year) anomalies of modelled $\Delta^{14}$C shown in panel b. d) Differences in the centennial $\Delta^{14}$C variations (panel c) from the control run. All model runs and panels are shown for the transition from preindustrial to perturbed conditions (transition 1, right), steady state of the perturbed conditions (steady state, middle), and the transition back to preindustrial carbon cycle conditions (transition 2, left). See also section 2.4.1.
Figure 5. Probability distributions for IntCal13-GICC05 timescale differences ($P_{scaled}(t_s)$, see section 2.1) for each 1,000-year window based on the mean of GRIP $^{10}$Be concentrations, fluxes, and their climate corrected versions (grey-scale patches in all panels). The gap in the GRIP $^{10}$Be record between 9,400 and 10,800 BP has been filled with data from the GISP2 ice core. Each probability distribution is centred on the mean age of the investigated window. a) Comparison to 95% probability intervals based on GRIP $^{10}$Be concentrations (solid orange), fluxes (solid blue) and their “climate corrected versions (dashed pink and green lines). b) Comparison to 95% confidence intervals based on the mean of GISP2 $^{10}$Be concentrations,
fluxes, and their climate corrected versions. Results for GISP2 are only shown for periods
where it has not been used to fill the gap in the GRIP record. c) Comparison to results based
on a different scaling (factors of 0.5 and 1.5 shown as blue and green lines, respectively) of
the GRIP $^{10}$Be record.
Figure 6. Rooted mean square error (RMSE$_{10^{\text{Be}}}$, see text) of synchronized centennial IntCal13 and $^{10}\text{Be}$-based $\Delta^{14}\text{C}$ variations as a function of different $^{10}\text{Be}$-scaling factors ($^{14}\text{C}:^{10}\text{Be}$ ratios). Results for the different versions of the GRIP$^{10}\text{Be}$ record are shown on the left, while GISP2 $^{10}\text{Be}$-based results are shown on the right.
Figure 7. Comparison of synchronized tree-ring (black) and ice core (grey) based $\Delta^{14}$C anomalies for $^{14}$C:$^{10}$Be ratios of 0.7 (top) and 1.1 (bottom).
Figure 8. Rooted mean square error (RMSE$_{10Be}$) of IntCal13 $\Delta^{14}C$ and $^{10}$Be based $\Delta^{14}C$ records from GRIP (left) and GISP2 (right) for different scalings of the $^{10}$Be based data after synchronization. The RMSE$_{10Be}$ has been calculated for binned data (bin size = 2.5 ‰, see text) taking IntCal $\Delta^{14}C$ errors into account.
Figure 9. IntCal13-GICC05 age transfer function (thick black line) and its 2σ confidence intervals (thin black lines) based on the probability distributions ($P_{scaled}(t_s)$, grey shading) obtained from comparing the GRIP $^{10}$Be-based $\Delta^{14}$C (mean of concentration, flux and climate corrections) and IntCal13 $\Delta^{14}$C records.
Figure 10. GRIP/GISP2 $^{10}\text{Be}$ based $\Delta^{14}\text{C}$ before (grey) and after (black) synchronization to IntCal13 (red) for the sections a) 3,500-4,500 years BP, b) 7,000-8,000 years BP, c) 10,000-11,000 years BP.
Figure 11. Comparison of the derived IntCal13-GICC05 timescale transfer function (black lines, this study) to the results by Muscheler et al. (2014, grey lines), and independent age markers that have been linked independently to the IntCal13 and GICC05 timescales at high precision (symbols). The results of this study and Muscheler et al. are shown with their respective 95% confidence intervals (dashed lines). The independent age markers are plotted as the difference between their estimated ages based on radiocarbon dating (Saksunarvatn Ash, Santorini), historical documents (Vesuvius) and dendrochronology (775 and 994 AD events), and their respective GICC05-ages. The plotted 1σ error bars largely reflect uncertainties in the radiocarbon-dating and calibration of the Saksunarvatn Ash (Lohne et al., 2013) and the Santorini eruption (Friedrich et al., 2006). Note that the identification of the Santorini tephra in ice cores has been challenged based on its geochemistry (Pearce et al., 2004).
Figure 12. Top: Comparison of the derived IntCal13-GICC05 transfer function (thin grey lines and shading, dashed lines denote the 95% confidence interval) to the GICC05 maximum counting error (bold grey lines). Bottom: Same as above but expressed as the rate of change (yrs/yr) of the GICC05 maximum counting error and the derived timescale transfer function.